

1 All points on the circle at $z = 0$ are the same distance $\sqrt{R^2 + z^2}$ away from $\vec{r} = z\vec{e}_z$, where we want to find the retarded potentials. The retardation condition (6.3.3) is then

$$t_{\text{ret}} + \frac{1}{c}\sqrt{R^2 + z^2} = t,$$

and since $\vec{R}(t_{\text{ret}})$ and $\vec{V}(t_{\text{ret}})$ are both perpendicular to \vec{r} , we have — for *this* geometry — the retarded potentials

$$\Phi(\vec{r}, t) = \frac{e}{\sqrt{R^2 + z^2}} \quad \text{for } \vec{r} = z\vec{e}_z$$

and

$$\vec{A}(\vec{r}, t) = \frac{1}{c}\vec{V}(t_{\text{ret}})\frac{e}{\sqrt{R^2 + z^2}} \quad \text{for } \vec{r} = z\vec{e}_z,$$

where

$$\vec{V}(t_{\text{ret}}) = v[-\vec{e}_x \sin \varphi + \vec{e}_y \cos \varphi]$$

with

$$\varphi = \frac{vt_{\text{ret}}}{R} = \frac{vt}{R} - \frac{v}{c}\sqrt{1 + (z/R)^2}.$$

2 The charge and current densities are

$$\begin{aligned} \rho(\vec{r}) &= \frac{e}{4\pi R^2}\delta(r - R), \\ \vec{j}(\vec{r}) &= \vec{\omega} \times \vec{r}\rho(\vec{r}, t) = \vec{\nabla} \times \left[\frac{e}{4\pi R}\vec{\omega}\eta(R - r) \right]. \end{aligned}$$

The electric field is

$$\vec{E}(\vec{r}) = \frac{\vec{r}}{r^3}4\pi \int_0^r dr' r'^2 \frac{e}{4\pi R^2}\delta(r' - R) = \eta(r - R)e\frac{\vec{r}}{r^3},$$

that is: no electric field inside the sphere, and outside it is the Coulomb field of the net charge.

The magnetic dipole moment is

$$\vec{\mu} = \frac{e}{8\pi Rc} \int (d\vec{r}) \underbrace{\vec{r} \times [\vec{\nabla} \times \vec{\omega}\eta(R - r)]}_{\rightarrow 2\vec{\omega}\eta(R - r) \text{ [see (3.6.8)]}} = \frac{e}{4\pi Rc}\vec{\omega}\frac{4\pi R^3}{3} = \frac{eR^2}{3c}\vec{\omega},$$

so that

$$\vec{j}(\vec{r}) = \vec{\nabla} \times \left[\frac{3c}{4\pi R^3}\vec{\mu}\eta(R - r) \right].$$

This gives first the vector potential

$$\begin{aligned}
\vec{A}(\vec{r}) &= \frac{1}{c} \int (d\vec{r}') \frac{\vec{j}(\vec{r}')}{|\vec{r} - \vec{r}'|} = -\frac{3}{4\pi R^3} \vec{\mu} \times \int (d\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} \vec{\nabla}' \eta(R - r') \\
&= -\frac{3}{4\pi R^3} \vec{\mu} \times \vec{\nabla} \int (d\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} \eta(R - r') \\
&= \frac{3}{4\pi R^3} \vec{\mu} \times \frac{\vec{r}}{r^3} 4\pi \int_0^r dr' r'^2 \eta(R - r') \\
&= \frac{\vec{\mu} \times \vec{r}}{R^3 r^3} \min\{R^3, r^3\} = \frac{\vec{\mu} \times \vec{r}}{\max\{R^3, r^3\}} \\
&= \eta(R - r) \frac{\vec{\mu} \times \vec{r}}{R^3} + \eta(r - R) \frac{\vec{\mu} \times \vec{r}}{r^3}
\end{aligned}$$

and then the magnetic field

$$\vec{B}(\vec{r}) = \vec{\nabla} \times \vec{A}(\vec{r}) = \eta(R - r) \frac{2\vec{\mu}}{R^3} + \eta(r - R) \frac{3\vec{\mu} \cdot \vec{r} \vec{r} - r^2 \vec{\mu}}{r^5},$$

that is: a constant magnetic field inside the sphere, and outside it is the field of the magnetic dipole $\vec{\mu}$.

3 In (see Exercise 33)

$$-\frac{dE}{dt} \Big|_{\text{rad}} = \frac{2e^2}{3c^3} \gamma^6 \left[\left(\frac{d\vec{v}}{dt} \right)^2 - \left(\frac{\vec{v}}{c} \times \frac{d\vec{v}}{dt} \right)^2 \right]$$

we insert what applies for motion on a circle with constant speed, that is:

$$\frac{d}{dt} \vec{v} = \vec{\omega}_0 \times \vec{v}, \quad v = \omega_0 R, \quad \vec{v} \times \frac{d}{dt} \vec{v} = v^2 \vec{\omega}_0,$$

and so get

$$-\frac{dE}{dt} \Big|_{\text{rad}} = \frac{2e^2}{3c^3} \gamma^6 \underbrace{\left[(\omega_0 v)^2 - (v^2 \omega_0 / c)^2 \right]}_{=(\omega_0 v)^2 / \gamma^2} = \frac{2e^2}{3c^3} \gamma^4 (\omega_0 v)^2 = \frac{2}{3} \omega_0 \frac{e^2}{R} \beta^3 \gamma^4,$$

which is the familiar result.