

PC5203 Advanced Solid State Physics

Weeks 7-9, due Monday 18 Oct 2021

[main concepts to cover: definitions of Green's functions, fluctuation-dissipation theorem, contour ordered Green's functions, Langreth rule, Caroli/Landauer formula]

1. (a) If the Hamiltonian is a single quantum dot of electron with $\hat{H} = \epsilon c^+ c$ in thermal equilibrium, show that the fluctuation-dissipation theorem is trivially true, i.e., the Green's functions in energy space E satisfy $g^<(E) = -f(E)(g^r(E) - g^a(E))$, here $f(E) = 1/(e^{\beta(E-\mu)} + 1)$ is the Fermi function. (b) Using the technique of diagonalizing an N -site Hamiltonian $\hat{H} = c^+ H c$, here the single particle H without the hat is an $N \times N$ Hermitian matrix, c is a column vector of N fermi annihilation operators, and c^+ is a row vector of Hermitian conjugate of c , show that the fluctuation-dissipation theorem still holds, i.e.,

$$G^< = -f(G^r - G^a)$$

here the Green's functions are matrices defined in the site space and in time domain t and Fourier transformed into energy domain E . (c) Using the Lehmann representation, show that the fluctuation-dissipation theorem holds in thermal equilibrium, for any Hamiltonian \hat{H} (with interaction or without) for the electron Green's functions.

2. Consider the time ordering operator T . (a) Define the precise meaning of time ordering operator on the product of the Hamiltonian operators: $T(H(t_1)H(t_2))$. (b) Use your definition to show that the following equation is true:

$$\frac{1}{2!} \int_{t'}^t dt_1 \int_{t'}^t dt_2 T(H(t_1)H(t_2)) = \int_{t'}^t dt_1 H(t_1) \int_{t'}^{t_1} dt_2 H(t_2)$$

where we assume $t > t'$. Note that the domain of integration on the left is on a square while the domain of the integral on the right-hand side is on a triangle. Draw illustrative pictures of the integration domains on the left and right side of the equal sign. (c) How would you modify the above equation if time order T is replaced by anti-time order \bar{T} ?

3. One form of the Langreth theorem says, if the convolution in contour time is $C(\tau, \tau') = \int A(\tau, \tau'')B(\tau'', \tau')d\tau''$, then, for time-translationally invariant system in energy domain it is

$$C^<(E) = A^r(E)B^<(E) + A^<(E)B^a(E).$$

Here A, B, C are matrices. Time-translationally invariant system means, e.g., $A(t, t') = A(t - t')$. Show this explicitly by elementary means, i.e., use the fact that $C^<(t) = C^{+-}(t, 0)$, $A^r = A^{++} - A^{+-}$, etc, and the meaning of contour integral in τ in relation to real time t , and the convolution theorem for Fourier transform.

4. Consider the electron transport in a single quantum dot with the center Hamiltonian $H_c = \epsilon c^\dagger c$, weakly coupled to two leads called left and right bath, as shown below in the figure. The interaction of the quantum dot with lead will be approximated by a retarded self energy (through the Dyson equation) which is constant and purely imaginary, such that the left lead $\Gamma_L = i(\Sigma_L^r - \Sigma_L^a) = -2 \text{Im} \Sigma_L^r$ is a constant independent of energy E , and similarly for the right lead. This is known as a wide-band approximation. (a) Determine the full Green's function $G^r(E)$, (b) Use the Caroli formula $\text{Tr}(G^r \Gamma_L G^a \Gamma_R)$ together with the Landauer formula to show that the current is given by the formula

$$I_L = \frac{-e}{\hbar} \frac{\Gamma_L \Gamma_R}{\Gamma_L + \Gamma_R} (f_L(\epsilon) - f_R(\epsilon)),$$

in the weak coupling limit, that is, you can assume Γ_L and Γ_R are small (thus a peak can be approximated as a delta function, and the integration in energy E can be performed approximately). Here f is the Fermi function for the left lead at β_L, μ_L or similarly the right lead at β_R, μ_R evaluated at the energy level ϵ of the quantum dot.

