

Question Test 2/1

Do not write on either margin

Write answers on this side of the paper only.

1 These are Gaussian wave functions so that they are minimum uncertainty wave functions. Therefore, the position wave function for ket  $|2\rangle$  is immediately available: It has the appearance of  $\langle x|1\rangle$  but with  $a$  replaced by  $\frac{\hbar}{2b}$ .

$$\langle x|2\rangle = \frac{(2\pi)^{-1/4}}{\sqrt{\frac{\hbar}{2b}}} e^{-\left(\frac{bx}{\hbar}\right)^2}$$

Alternatively, you can get this by Fourier transforming  $\langle p|2\rangle$ . Then

$$\begin{aligned} \langle 1|2\rangle &= \int dx \langle 1|x\rangle \langle x|2\rangle \\ &= \frac{(2\pi)^{-1/2}}{\sqrt{\frac{\hbar a}{2b}}} \int dx e^{-\left(\frac{x}{2a}\right)^2 - \left(\frac{bx}{\hbar}\right)^2} \\ &= \sqrt{\frac{2b}{2\pi\hbar a}} \sqrt{\frac{\pi}{\left(\frac{1}{2a}\right)^2 + \left(\frac{b}{\hbar}\right)^2}} = \sqrt{\frac{b(2\hbar a)^2}{\hbar a \hbar^2 + (2ab)^2}} \\ &= \sqrt{\frac{4\hbar ab}{\hbar^2 + (2ab)^2}} \end{aligned}$$

$$\text{and } |\langle 1|2\rangle|^2 = \frac{4ab/\hbar}{1 + (2ab/\hbar)^2}$$

For  $a = \frac{\hbar}{2b}$  or  $2ab/\hbar = 1$ , we have  $\langle 1|2\rangle = 1$ , and get  $|\langle 1|2\rangle|^2 = 1$ , as we should.

$$2(a) \frac{d}{dt} P = -\frac{\partial H}{\partial X} = 2\Omega P \sim P(t) = P(t_0) e^{2\Omega T}$$

$$\frac{d}{dt} X = \frac{\partial H}{\partial P} = -2\Omega X \sim X(t) = X(t_0) e^{-2\Omega T}$$

$$(b) [X(t), P(t_0)] = e^{-2\Omega T} [X(t_0), P(t_0)] = i\hbar e^{-2\Omega T}$$

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$$(c) i\hbar \frac{\partial}{\partial t} \langle x, t | p, t_0 \rangle = \langle x, t | H | p, t_0 \rangle$$

with  $H$  at time  $t$ :

$$\begin{aligned} H &= -\Omega (X(t)P(t) + P(t)X(t)) \\ &= -2\Omega X(t)P(t) + i\hbar\Omega \\ &= -2\Omega X(t) e^{2\Omega T} P(t_0) + i\hbar\Omega \end{aligned}$$

so that

$$i\hbar \frac{\partial}{\partial t} \langle x, t | p, t_0 \rangle = (-2\Omega X(t) e^{2\Omega T} + i\hbar\Omega) \langle x, t | p, t_0 \rangle$$

or

$$\begin{aligned} \frac{\partial}{\partial t} \log \langle x, t | p, t_0 \rangle &= i \frac{XP}{\hbar} 2\Omega e^{2\Omega T} + \Omega \\ &= \frac{\partial}{\partial t} \left( i \frac{XP}{\hbar} e^{2\Omega T} + \Omega T \right) \end{aligned}$$

It follows that

$$\langle x, t | p, t_0 \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{i \frac{XP}{\hbar} e^{2\Omega T}} e^{-\Omega T}$$

after incorporating the initial condition

$$\langle x, t | p, t_0 \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ixp/\hbar} \text{ for } T \rightarrow 0.$$

Now,

$$\begin{aligned} \langle x, t | x', t_0 \rangle &= \int dp \langle x, t | p, t_0 \rangle \langle p, t_0 | x', t_0 \rangle \\ &= \int dp \frac{1}{\sqrt{2\pi\hbar}} e^{ixpe^{2\Omega T}/\hbar} e^{-\Omega T} \frac{e^{-ipx'/\hbar}}{\sqrt{2\pi\hbar}} \\ &= e^{-\Omega T} \delta(xe^{2\Omega T} - x') \\ &= \delta(xe^{-\Omega T} - x'e^{-\Omega T}). \end{aligned}$$

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$$\begin{aligned} \text{[3] We have } \rho(A^+(t), A(t), t) &= \rho(A^+(t_0), A(t_0), t_0) \\ &= e^{-\frac{1}{2}(A^+(t_0) - \alpha^*), (A(t_0) - \alpha)} \end{aligned}$$

$$\text{and } A^+(t) = e^{i\omega t} A^+(t_0), \quad A(t) = e^{-i\omega t} A(t_0)$$

$$\text{or } A^+(t_0) = e^{-i\omega t} A^+(t), \quad A(t_0) = e^{i\omega t} A(t)$$

$$\begin{aligned} \text{So that} \\ \rho(A^+, A, t) &= e^{-\frac{1}{2}(A^+ e^{-i\omega t} - \alpha^*), (A e^{i\omega t} - \alpha)} \\ &= e^{-\frac{1}{2}(A^+ - \alpha^* e^{i\omega t}), (A - \alpha e^{-i\omega t})} \end{aligned}$$

$$\begin{aligned} \text{[4] (a) With } |1\rangle &= |0\rangle \frac{1}{\sqrt{2}}, \quad \langle 1| A^+ = \frac{1}{\sqrt{2}} \langle 0| \\ A^2 |1\rangle &= 0, \quad \langle 1| A^{+2} = 0, \end{aligned}$$

$$\text{we get } \langle A \rangle = \langle |0\rangle \frac{1}{\sqrt{2}} = \frac{1}{2}, \quad \langle A^+ \rangle = \frac{1}{2}$$

$$\langle A^2 \rangle = 0, \quad \langle A^{+2} \rangle = 0$$

$$\text{(b) Also will need } \langle A^+ A \rangle = \frac{1}{2}, \quad \langle A A^+ \rangle = \frac{3}{2}$$

$$\text{in } \langle X^2 \rangle = \left\langle \frac{\ell^2}{2} (A^{+2} + A^+ A + A A^+ + A^2) \right\rangle = \ell^2,$$

$$\langle P^2 \rangle = \left\langle \frac{\hbar^2}{2e^2} (-A^{+2} + A^+ A + A A^+ - A^2) \right\rangle = \left(\frac{\hbar}{e}\right)^2$$

$$\text{With } \langle X \rangle = \frac{\ell}{\sqrt{2}} \langle (A^+ + A) \rangle = \frac{\ell}{\sqrt{2}}$$

$$\text{and } \langle P \rangle = \frac{\hbar}{\sqrt{2}e} \langle (iA^+ - iA) \rangle = 0$$

this gives

$$\left. \begin{aligned} \Delta X &= \sqrt{\ell^2 - \left(\frac{\ell}{\sqrt{2}}\right)^2} = \frac{\ell}{\sqrt{2}} \\ \Delta P &= \sqrt{\left(\frac{\hbar}{e}\right)^2 - 0^2} = \frac{\hbar}{e} \end{aligned} \right\} \Delta X \Delta P = \frac{\hbar}{\sqrt{2}} > \frac{\hbar}{2}$$

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5(a) Use the completeness relation for coherent states,

$$\begin{aligned} \text{tr}\{F\} &= \text{tr}\left\{F \int \frac{dx dp}{2\pi\hbar} \frac{|a'\rangle\langle a^*|}{\langle a^*|a'\rangle}\right\} \\ &= \int \frac{dx dp}{2\pi\hbar} \frac{\langle a'|F|a^*\rangle}{\langle a^*|a'\rangle} \\ &= \int \frac{dx dp}{2\pi\hbar} f(a^*, a') \end{aligned}$$

chosen  
parameterization

and take the parameterization  $a' = \frac{x}{\epsilon}$ ,  $a^* = -i\frac{p}{\hbar}$   
to get

$$\text{tr}\{F\} = \int \frac{dx dp}{2\pi\hbar} f\left(-i\frac{p}{\hbar}, \frac{x}{\epsilon}\right).$$

$$\begin{aligned} \text{(b)} \quad \text{tr}\{e^{-\lambda A^+}; A\} &= \int \frac{dx dp}{2\pi\hbar} e^{-\lambda\left(\frac{i p}{\hbar}\right) \frac{x}{\epsilon}} \\ &= \int \frac{dx dp}{2\pi\hbar} e^{i\lambda x p / \hbar} = \int dx \delta(\lambda x) = \frac{1}{\lambda}. \end{aligned}$$