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$$\begin{aligned}
 \text{[1(a)} \quad \text{tr}\{F^{\dagger}G\} &= \sum_{k,e} \sum_{k',e'} f_{ke}^* g_{k'e'} \text{tr}\{V^{-e} U^{-k} U^{k'} V^{e'}\} \\
 &= \sum_{k,e} \sum_{k',e'} f_{ke}^* g_{k'e'} N \delta_{kk'} \delta_{ee'} \\
 &= N \sum_{k,e} f_{ke}^* g_{ke} .
 \end{aligned}$$

$$\begin{aligned}
 \text{(b)} \quad \sum_{k=1}^N \langle u_k | &= \sum_{k,e=1}^N \langle u_k | v_e \rangle \langle v_e | \\
 &= \frac{1}{\sqrt{N}} \sum_{k,e=1}^N e^{i \frac{2\pi}{N} k e} \langle v_e | \\
 &= \frac{1}{\sqrt{N}} \sum_{e=1}^N N \delta_{eN} = \sqrt{N} \langle v_N | .
 \end{aligned}$$

Likewise, $\sum_{e=1}^N |v_e\rangle = |u_N\rangle \sqrt{N}$.

[2(a) Characteristically $|\langle q | \gamma \rangle|^2 = \frac{1}{2\pi}$ does not depend on the quantum numbers q and γ , so that Q, Γ are a complementary pair.

$$\begin{aligned}
 \text{(b)} \quad \int_0^{\infty} \frac{dq}{q} \langle \gamma | q \rangle \langle q | \gamma' \rangle &= \int_0^{\infty} \frac{dq}{q} q^{-i(\gamma-\gamma')/2\pi} \\
 \text{q} = e^x &\Rightarrow \int_{-\infty}^{\infty} dx \frac{1}{2\pi} e^{-ix(\gamma-\gamma')} = \delta(\gamma-\gamma') ,
 \end{aligned}$$

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which states that

$$\delta(x-x') = \langle x|x' \rangle = \langle x | \left(\int_0^\infty \frac{dq}{q} |q\rangle \langle q| \right) |x' \rangle$$

$$\text{and implies } \int_0^\infty \frac{dq}{q} |q\rangle \langle q| = 1.$$

$$\begin{aligned} \int_{-\infty}^{\infty} dy \langle q|x \rangle \langle x|q' \rangle &= \int_{-\infty}^{\infty} dy \left(\frac{q}{q'} \right)^{ix} / 2\pi \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dy e^{ix \log(q/q')} = \delta(\log(q/q')) \\ &= q \delta(q-q') \\ &= \langle q|q' \rangle, \end{aligned}$$

$$\text{and } \int_{-\infty}^{\infty} dy |x\rangle \langle x| = 1 \text{ follows.}$$

$$\begin{aligned} \text{(c) Consider } \langle q|Q^{ix'}|x \rangle &= q^{ix'} \langle q|x \rangle = \frac{q^{i(x+x')}}{\sqrt{2\pi}} \\ &= \langle q|x+x' \rangle, \end{aligned}$$

and the completeness of the bras $\langle q|$
implies

$$Q^{ix'}|x \rangle = |x+x' \rangle.$$

$$\begin{aligned} \text{(d) Compare } e^{i\beta\Gamma} Q^{ix} |x' \rangle &= e^{i\beta\Gamma} |x+x' \rangle = |x+x' \rangle e^{i\beta(x+x')} \\ \text{with } Q^{ix} e^{i\beta\Gamma} |x' \rangle &= Q^{ix} |x' \rangle e^{i\beta x'} = |x+x' \rangle e^{i\beta x'} \\ \text{and conclude that } e^{i\beta\Gamma} Q^{ix} &= e^{i\beta x} Q^{ix} e^{i\beta\Gamma}. \end{aligned}$$

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[3] (a) For example, consider an arbitrary power of A :

$$e^{-\epsilon B} A^n e^{\epsilon B} = (e^{-\epsilon B} A e^{\epsilon B})^n$$

by inspection, then

$$e^{-\epsilon B} e^A e^{\epsilon B} = e^{e^{-\epsilon B} A e^{\epsilon B}}$$

follows by summation of the power series for e^A .

(b) To first order in ϵ , we have

$$e^{-\epsilon B} e^A e^{\epsilon B} = e^A + \epsilon [e^A, B]$$

and

$$e^{-\epsilon B} A e^{\epsilon B} = A + \underbrace{\epsilon [A, B]}_{\equiv \delta A} = A + \delta A.$$

$$\begin{aligned} \text{So } \delta e^A &= \epsilon [e^A, B] = \int_0^1 dx e^{(1-x)A} \delta A e^{xA} \\ &= \int_0^1 dx e^{(1-x)A} \epsilon [A, B] e^{xA}, \end{aligned}$$

$$\text{giving } [e^A, B] = \int_0^1 dx e^{(1-x)A} [A, B] e^{xA}.$$

[4] (a) For any eigenvalue $a > 0$ of A we have

$$\int_0^\infty d\alpha \left(\frac{1}{\alpha+1} - \frac{1}{\alpha+a} \right) = \int_0^\infty d\alpha \frac{d}{d\alpha} \log \frac{\alpha+1}{\alpha+a}$$

$$= \log \frac{\alpha+1}{\alpha+a} \Big|_{\alpha=0}^{\infty} = \log 1 - \log \frac{1}{a} = \log a$$

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and

$$\int_0^{\infty} d\beta \frac{e^{-\beta} - e^{-\beta a}}{\beta} = \int_0^{\infty} d\beta \int_1^a d\alpha e^{-\alpha\beta}$$

$$= \int_1^a d\alpha \int_0^{\infty} d\beta e^{-\alpha\beta} = \int_1^a d\alpha \frac{1}{\alpha} = \log a,$$

so that the two integrals give the correct values for all eigenvalues of a , as they should.

(b) Recall $\delta \frac{1}{x} = -\frac{1}{x^2} \delta x$ and use

this for $x = \alpha + A$, $\delta x = \delta A$, then

$$\delta \log A = \delta \int_0^{\infty} d\alpha \left(\frac{1}{\alpha+1} - \frac{1}{\alpha+A} \right)$$

$$= \int_0^{\infty} d\alpha \frac{1}{\alpha+A} \delta A \frac{1}{\alpha+A}.$$