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1 (a) From Test 2 / Problem 2 we know that these relations hold for the two single delta potentials, with

$$\begin{pmatrix} t & r \\ r & t \end{pmatrix} = e^{i\alpha} \begin{pmatrix} \frac{ka}{ka-i} & \frac{i}{ka-i} \\ \frac{i}{ka-i} & \frac{ka}{ka-i} \end{pmatrix} \text{ where } \alpha = kL.$$

Writing $ka = \cot\beta$, we have therefore

$$t = e^{i\alpha} \frac{\cot\beta}{\cot\beta - i} = e^{i\alpha} \frac{\cos\beta}{\cos\beta - i\sin\beta} = e^{i(\alpha+\beta)} \cos\beta$$

and

$$r = \frac{i}{ka} t = i t \tan\beta = i e^{i(\alpha+\beta)} \sin\beta.$$

(b) We have the two equations

$$\begin{pmatrix} \phi_+(L) \\ \phi_-(-L) \end{pmatrix} = \begin{pmatrix} 0 & r \\ r & 0 \end{pmatrix} \begin{pmatrix} \phi_+(-L) \\ \phi_-(L) \end{pmatrix} + t \begin{pmatrix} \phi_+(0) \\ \phi_-(0) \end{pmatrix}$$

where we need $\phi_{\pm}(0)$, which we get from the other two equations, namely

$$\begin{pmatrix} 1 & -r \\ -r & 1 \end{pmatrix} \begin{pmatrix} \phi_+(0) \\ \phi_-(0) \end{pmatrix} = t \begin{pmatrix} \phi_+(-L) \\ \phi_-(L) \end{pmatrix}.$$

Accordingly,

$$\begin{aligned} \begin{pmatrix} TR \\ RT \end{pmatrix} &= \begin{pmatrix} 0 & r \\ r & 0 \end{pmatrix} + t^2 \begin{pmatrix} 1 & -r \\ -r & 1 \end{pmatrix}^{-1} \\ &= \begin{pmatrix} 0 & r \\ r & 0 \end{pmatrix} + \frac{t^2}{1-r^2} \begin{pmatrix} 1 & r \\ r & 1 \end{pmatrix}, \end{aligned}$$

so that $T = \frac{t^2}{1-r^2}$

and $R = \frac{r}{1-r^2} (1-r^2+t^2).$

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(c) We need $0 = 1 - r^2 + t^2 = 1 + e^{2i(\alpha+\beta)}$
or

$$\cos(\alpha+\beta) = 0,$$

or $\cos\alpha \cos\beta = \sin\alpha \sin\beta$,
so that

$$\tan\alpha = \cot\beta$$

or, finally

$$\boxed{\tan(kL) = ka}.$$

[2] (a) For the Yukawa potential we have

$$-\frac{M}{2\pi\hbar^2} \int (d\vec{r}') e^{i(\vec{k}-\vec{k}')\cdot\vec{r}} \frac{V_0}{kr} e^{-kr}$$

$$= -\frac{2MV_0/k}{\hbar^2(k^2+q^2)} \equiv f_0(\theta) \text{ with } q = 2k \sin \frac{\theta}{2}.$$

Therefore,

$$f(\theta) = 2 \cos((\vec{k}-\vec{k}')\cdot\vec{a}) f_0(\theta)$$

[see Test 2/Problem 4] where $\vec{k}\cdot\vec{a} = ka$
and $\vec{k}'\cdot\vec{a} = \vec{k}'\cdot\vec{k} a/k = ka \cos\theta$,
so that

$$f(\theta) = 2 \cos(ka(1-\cos\theta)) f_0(\theta)$$

$$= 2 \cos(2ka(\sin \frac{\theta}{2})^2) f_0(\theta),$$

and $\frac{dS}{d\Omega} = 4 \left(\cos(2ka(\sin \frac{\theta}{2})^2) \right)^2 |f_0(\theta)|^2.$

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(b) For $\frac{d\delta}{d\Omega} = 0$ we need $\cos(ka(1-\cos\theta)) = 0$,
so that

$$\cos\theta = \frac{2}{3}: \quad \cos\left(\frac{1}{3}ka\right) = 0,$$

$$\cos\theta = 0: \quad \cos(ka) = 0,$$

$$\cos\theta = -\frac{2}{3}: \quad \cos\left(\frac{5}{3}ka\right) = 0,$$

which tell us that $ka = \frac{3\pi}{2} = 2\pi \times \frac{3}{4}$.
It follows that

$$\boxed{a = \frac{3}{4}\lambda}$$

3 (a) We have

$$\begin{aligned} (A^+\sigma + \sigma^+A)^2 &= A^+A\sigma\sigma^+ + AA^+\sigma^+\sigma \\ &= A^+A + \underbrace{\sigma\sigma^+ + \sigma^+\sigma}_{= 1 - \sigma^+\sigma} \\ &= A^+A + \sigma^+\sigma, \end{aligned}$$

Since $\sigma^2 = 0$ and $\sigma^{+2} = 0$. This tells us that, essentially, H_0 is the square of H_1 , so that the two parts of H commute, implying

$$[H(t_1), H(t_2)] = 0$$

for all t_1 and t_2 .

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$$(b) \alpha(t) = \langle e, 0, t | \rangle, \beta(t) = \langle g, 1, t | 0 \rangle \text{ give}$$

$$i\hbar \frac{\partial}{\partial t} \alpha(t) = \langle e, 0, t | H | \rangle = \hbar\omega \langle e, 0, t | \rangle - \hbar\Omega \langle g, 1, t | \rangle$$

$$= \hbar\omega \alpha(t) - \hbar\Omega \beta(t),$$

$$i\hbar \frac{\partial}{\partial t} \beta(t) = \langle g, 1, t | H | \rangle = \hbar\omega \langle g, 1, t | \rangle - \hbar\omega \langle e, 0, t | \rangle$$

$$= \hbar\omega \beta(t) - \hbar\Omega \alpha(t),$$

$$\text{or } \frac{\partial}{\partial t} \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix} = i \begin{pmatrix} -\omega & \Omega(t) \\ \Omega(t) & -\omega \end{pmatrix} \begin{pmatrix} \alpha(t) \\ \beta(t) \end{pmatrix}.$$

(c) Since the eigenvalues of \uparrow are $(1, 1)$ and $(1, -1)$, the time dependence of $\alpha \pm \beta$ is particularly simple:

$$\frac{\partial}{\partial t} (\alpha(t) \pm \beta(t)) = i(-\omega \pm \Omega(t)) (\alpha(t) \pm \beta(t)),$$

with the consequence

$$\alpha(t) \pm \beta(t) = (\alpha(0) \pm \beta(0)) e^{-i\omega t \pm i \int_0^t \Omega(t') dt'}$$

For $\alpha(0) = 1, \beta(0) = 0$ and $t = T$, this gives

$$\alpha(T) \pm \beta(T) = e^{-i\omega T} e^{\pm i \int_0^T dt \Omega(t)}$$

or with

$$\int_0^T dt \Omega(t) = 2 \int_0^{T/2} dt \frac{2\pi t}{T^2} = \frac{\pi}{2},$$

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$$\alpha(T) \pm \beta(T) = e^{-i\omega T} e^{\pm i\frac{\pi}{2}} = \pm i e^{-i\omega T}$$

Thus, $\alpha(T) = 0$, $\beta(T) = i e^{-i\omega T}$ and the probability in question is

$$|\beta(T)|^2 = 1.$$

4 (a) One electron: $|1\uparrow\rangle$, the other: $|2\downarrow\rangle$, together $| \rangle = \left(\underbrace{|1\uparrow, 2\downarrow\rangle}_{\text{electron 1}} - \underbrace{|2\downarrow, 1\uparrow\rangle}_{\text{electron 2}} \right) / \sqrt{2}$.

Singlet component: $(|1\uparrow\downarrow\rangle - |1\downarrow\uparrow\rangle) / \sqrt{2}$, has spatial wave function

$$\begin{aligned} \psi_s(\vec{r}_1, \vec{r}_2) &= \langle \vec{r}_1, \vec{r}_2 | (|112\rangle + |211\rangle) / \sqrt{2} \rangle \\ &= \frac{1}{2} [\psi_1(\vec{r}_1) \psi_2(\vec{r}_2) + \psi_2(\vec{r}_1) \psi_1(\vec{r}_2)]; \end{aligned}$$

triplet components: $(|1\uparrow\downarrow\rangle + |1\downarrow\uparrow\rangle) / \sqrt{2}$, has spatial wave function

$$\begin{aligned} \psi_t(\vec{r}_1, \vec{r}_2) &= \langle \vec{r}_1, \vec{r}_2 | (|112\rangle - |211\rangle) / \sqrt{2} \rangle \\ &= \frac{1}{2} [\psi_1(\vec{r}_1) \psi_2(\vec{r}_2) - \psi_2(\vec{r}_1) \psi_1(\vec{r}_2)]. \end{aligned}$$

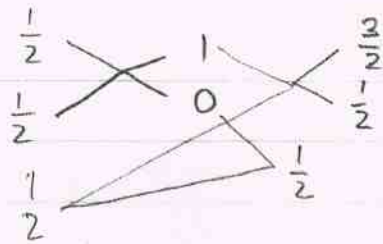
(b) These probabilities are

$$\begin{aligned} \int d\vec{r}_1 d\vec{r}_2 |\psi_s(\vec{r}_1, \vec{r}_2)|^2 &= \frac{1}{4} (1+1 + \langle 112 | \langle 211 | \\ &\quad + \langle 211 | \langle 112 |) \\ &= \frac{1}{2} + \frac{1}{2} |\langle 112 \rangle|^2 = \frac{1+8}{2} \end{aligned}$$

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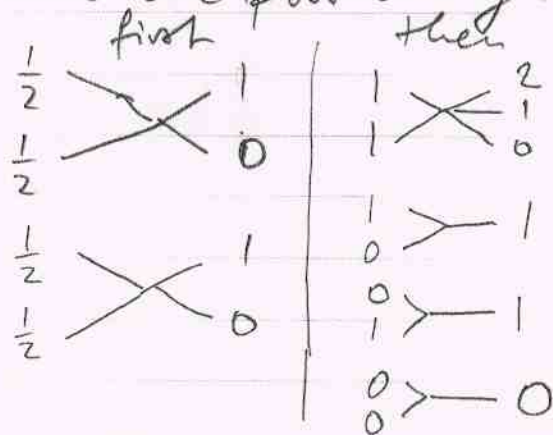
and $(d\vec{r}_1)(d\vec{r}_2) + \frac{1}{4} (\vec{r}_1, \vec{r}_2)^2 = \frac{1-\gamma}{2}$,
 respectively.

(c) The three angular momenta of $\frac{1}{2}$ each can be coupled in this manner



We begin with $2^3 = 8$ states which are grouped into 4 states with $s = \frac{3}{2}$ and 2 times 2 states with $s = \frac{1}{2}$, together $4 + 2 + 2 = 8$ states.

(d) Here one possibility is



so that the $2^4 = 16$ states are grouped into 5 states with $s = 2$, 3 times 3 states with $s = 1$, and 2 times 1 state with $s = 0$, together $5 + 3 \times 3 + 2 \times 1 = 16$ states.