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1 We use $[x, P] = i\hbar$ twice,

$$\begin{aligned} P X^3 P - X P X P X &= P X^3 P - (P X + i\hbar) X (X P - i\hbar) \\ &= P X^3 P - (P X^3 P - i\hbar P X^2 - i\hbar X^2 P + \hbar^2 X) \\ &= -i\hbar \underbrace{[X^2, P]}_{2i\hbar X} - \hbar^2 X = \underline{\hbar^2 X}. \end{aligned}$$

2 We have $f(x)^+ = \int dx (|x\rangle f(x) \langle x|)^+ = \int dx |x\rangle f(x)^* \langle x|$,
so that(a) $f(x) = f(x)^+$ requires $f(x) = f(x)^*$, that is: $f(x)$ must be real for all x values;(b) $f(x)^+ f(x) = 1$ requires $f(x)^* f(x) = 1 |f(x)|^2 = 1$,
that is: $f(x)$ must be a phase factor for all x values.3 Since $e^{ipx/\hbar} e^{ixP/\hbar} = e^{-ixp/\hbar} e^{ixP/\hbar} e^{ipx/\hbar}$, we
need $e^{-ixp/\hbar} = 1$, that is xp must be an
integer multiple of $2\pi\hbar$: $xp = 0, \pm 2\pi\hbar, \pm 4\pi\hbar, \dots$.4(a) To show that $UU^+ = 1$ we consider $\langle x | UU^+ | x' \rangle$,

$$\begin{aligned} \langle x | UU^+ | x' \rangle &= \int dx'' \langle x | U | x'' \rangle \langle x'' | U^+ | x' \rangle \\ &= \int dx'' \langle x | U | x'' \rangle (\langle x' | U | x'' \rangle)^* \end{aligned}$$

$$\begin{aligned} &= \int dx'' \frac{e^{ixx''/\alpha^2}}{\sqrt{2\pi}\alpha} \frac{e^{-ix'x''/\alpha^2}}{\sqrt{2\pi}\alpha} \\ &= \frac{1}{2\pi\alpha^2} \int dx'' e^{i(x-x')x''/\alpha^2} \end{aligned}$$

$$= \frac{1}{2\pi\alpha^2} 2\pi \delta\left(\frac{x-x'}{\alpha^2}\right) = \delta(x-x'),$$

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so that $\langle x | UU^\dagger | x' \rangle = \langle x | 1 | x' \rangle$ for all $|x\rangle$ and all $|x'\rangle$, implying $UU^\dagger = 1$ indeed.

4(b) We have

$$\begin{aligned} \langle x | U | p' \rangle &= \int dx' \langle x | U | x' \rangle \langle x' | p' \rangle \\ &= \int dx' \frac{e^{ixx'/\hbar}}{\sqrt{2\pi/\hbar}} \frac{e^{-ix'p'/\hbar}}{\sqrt{2\pi/\hbar}} \\ &= \frac{1}{2\pi a\sqrt{\hbar}} 2\pi \delta\left(\frac{x}{a^2} + \frac{p'}{\hbar}\right) \\ &= \frac{1}{\sqrt{\hbar}} \delta(x/a + ap'/\hbar). \end{aligned}$$

4(c) Now we get

$$\begin{aligned} \langle p | U | p' \rangle &= \int dx \langle p | x \rangle \langle x | U | p' \rangle \\ &= \int dx \frac{e^{-ipx/\hbar}}{\sqrt{2\pi/\hbar}} \frac{1}{\sqrt{\hbar}} \delta\left(\frac{x}{a} + \frac{ap'}{\hbar}\right) \\ &= \frac{a}{\sqrt{2\pi/\hbar}} e^{-ip(-\frac{a^2 p'}{\hbar})/\hbar} \\ &= \frac{a/\hbar}{\sqrt{2\pi}} e^{ia^2 pp'/\hbar^2}. \end{aligned}$$

4(d) We have

$$\begin{aligned} \langle x | U^2 | x' \rangle &= \int dx'' \langle x | U | x'' \rangle \langle x'' | U | x' \rangle \\ &= \int dx'' \frac{1}{2\pi a^2} e^{i(x-x''+x''x')/\hbar^2} \\ &= \frac{1}{2\pi a^2} 2\pi \delta\left(\frac{x+x'}{a^2}\right) = \delta(x+x') \\ &= \langle -x | x' \rangle \text{ for all } |x'\rangle, \end{aligned}$$

so that

$$\langle x | U^2 = \langle -x | \text{ for all } |x\rangle.$$

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$$\begin{aligned}
 5(a) \quad & \langle e^{ikx} \rangle = \int dx \langle |x\rangle e^{ikx} \langle x|1\rangle \\
 & = \int dx e^{ikx} |4(x)|^2 \\
 & = \int dx e^{-ikx} \underbrace{k e^{-2k|x|}}_{\substack{= \cos(kx) + i \sin(kx) \\ \text{even} \quad \text{odd, does not contribute}}} \\
 & = 2k \int_0^\infty e^{-2k|x|} \underbrace{\cos(kx)}_{\substack{= \operatorname{Re} e^{ikx}}} \\
 & = 2k \operatorname{Re} \frac{1}{2k - ik} = \underline{\underline{\frac{(2k)^2}{(2k)^2 + k^2}}}.
 \end{aligned}$$

$$\begin{aligned}
 5(b) \quad & \text{It is } \langle x|e^{\frac{i}{2}iaP/t}|1\rangle = \langle x + \frac{a}{2}|1\rangle = 4(x + \frac{a}{2}) \\
 & \text{and } \langle 1|e^{\frac{i}{2}iaP/t}|x\rangle = \langle 1|x - \frac{a}{2}\rangle = 4(x - \frac{a}{2}),
 \end{aligned}$$

so that

$$\begin{aligned}
 \langle e^{iaP/t} \rangle & = \int dx \langle 1|e^{\frac{i}{2}iaP/t}|x\rangle \langle x|e^{\frac{i}{2}iaP/t}|1\rangle \\
 & = \int dx 4(x - \frac{a}{2})^* 4(x + \frac{a}{2}) \\
 & = K \int dx e^{-K(|x - \frac{a}{2}| + |x + \frac{a}{2}|)}.
 \end{aligned}$$

Since $|x - \frac{a}{2}| + |x + \frac{a}{2}|$ is even in both x and a , we can take twice the integral $\int_0^\infty dx$ and replace a by $|a|$. Then

$$\langle e^{iaP/t} \rangle = 2K \int_0^\infty dx e^{-K(|x - \frac{|a|}{2}| + |x + \frac{|a|}{2}|)}$$

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$$= 2K \int_0^{|a|/2} dx e^{-k|x|} + 2K \int_{|a|/2}^{\infty} dx e^{-2kx}$$

$$= k|a| e^{-k|a|} + e^{-k|a|} = \underline{(1+k|a|) e^{-k|a|}}.$$

5(c) Upon expanding $\langle e^{ikX} \rangle = \frac{(2k)^2}{(2k)^2 + h^2} = 1 - \frac{h^2}{(2k)^2 + h^2}$
in powers of k up to k^2 ,

$$\langle (1 + ikX - \frac{1}{2}k^2 X^2) \rangle = 1 - \frac{h^2}{(2k)^2},$$

we read off that $\langle X \rangle = 0$, $\langle X^2 \rangle = \frac{1}{2k^2}$.
Likewise

$$\langle e^{iaP/\hbar} \rangle = (1 + k|a|) \left(1 - k|a| + \frac{1}{2}(ka)^2 + \dots \right)$$

gives $\langle (1 + i\frac{a}{\hbar}P - \frac{1}{2}(\frac{a}{\hbar}P)^2) \rangle = 1 - \frac{1}{2}(ka)^2$,

which yields

$$\langle P \rangle = 0, \quad \langle P^2 \rangle = (\hbar k)^2.$$

The spreads $\delta X = \frac{1}{\sqrt{2}k}$, $\delta P = \hbar k$ follow immediately, and their product obeys the uncertainty relation

$$\delta X \delta P = \frac{\hbar}{\sqrt{2}} = \sqrt{2} \frac{\hbar}{2} > \frac{\hbar}{2},$$

as it should.