

## Question. Text 2.1.1

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1 We have

$$\frac{d}{dt} P(t) = 0, \quad P(t) = P(t_0)$$

$$\text{and } \frac{d}{dt} X(t) = \frac{1}{M(t)} P(t) = \frac{1}{M(t)} P(t_0),$$

$$\text{so that } X(t) = X(t_0) + P(t_0) \int_{t_0}^t \frac{dt'}{M(t')}.$$

It follows that

$$\begin{aligned} i\hbar \frac{\partial}{\partial t} \langle x, t | p, t_0 \rangle &= \langle x, t | H(t) | p, t_0 \rangle \\ &= \langle x, t | \frac{1}{2M(t)} P(t_0)^2 | p, t_0 \rangle \\ &= \frac{p^2}{2M(t)} \langle x, t | p, t_0 \rangle; \end{aligned}$$

$$\begin{aligned} -i\hbar \frac{\partial}{\partial t_0} \langle x, t | p, t_0 \rangle &= \langle x, t | H(t_0) | p, t_0 \rangle \\ &= \langle x, t | \frac{1}{2M(t_0)} P(t_0)^2 | p, t_0 \rangle \\ &= \frac{p^2}{2M(t_0)} \langle x, t | p, t_0 \rangle; \end{aligned}$$

$$\begin{aligned} \frac{\hbar}{i} \frac{\partial}{\partial x} \langle x, t | p, t_0 \rangle &= \langle x, t | P(t) | p, t_0 \rangle \\ &= \langle x, t | P(t_0) | p, t_0 \rangle \\ &= p \langle x, t | p, t_0 \rangle; \end{aligned}$$

$$\begin{aligned} \frac{\hbar}{i} \frac{\partial}{\partial p} \langle x, t | p, t_0 \rangle &= \langle x, t | X(t_0) | p, t_0 \rangle \\ &= \langle x, t | (X(t) - P(t) \int_{t_0}^t \frac{dt'}{M(t')}) | p, t_0 \rangle \\ &= (x - p \int_{t_0}^t \frac{dt'}{M(t')}) \langle x, t | p, t_0 \rangle. \end{aligned}$$

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Taken together these differential equations imply

$$\langle x, t | p, t_0 \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{i x p / \hbar} e^{-\frac{i}{\hbar} \frac{p^2}{2} \int_{t_0}^t dt' M(t')}$$

with the prefactor determined by the  $t \rightarrow t_0$  limit, or the familiar  $M(t) = \text{const}$  result.

[2] We have  $V^\dagger = \sum_{n=0}^{\infty} |n\rangle \langle n+1|$ , so that

$$V^\dagger V = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} |m\rangle \underbrace{\langle m+1 | n+1 \rangle}_{\delta_{n,m}} \langle n| = \sum_{n=0}^{\infty} |n\rangle \langle n| = 1, \text{ indeed.}$$

$$\begin{aligned} \text{But } V V^\dagger &= \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} |n+1\rangle \underbrace{\langle n | m \rangle}_{\delta_{nm}} \langle m+1| \\ &= \sum_{n=0}^{\infty} |n+1\rangle \langle n+1| = \sum_{n=1}^{\infty} |n\rangle \langle n| \\ &= 1 - |0\rangle \langle 0| \neq 1. \end{aligned}$$

In particular,  $\langle n=0 | V = 0$ , which is impossible if  $V$  were unitary.

$$\begin{aligned} \text{Further, } V &= A^\dagger \sum_{n=0}^{\infty} |n\rangle \frac{1}{\sqrt{n+1}} \langle n| \\ &= A^\dagger \frac{1}{\sqrt{A^\dagger A + 1}} \end{aligned}$$

$$\text{so that } W(A^\dagger A) = (A^\dagger A + 1)^{-1/2}.$$

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$$\begin{aligned} \boxed{3} \text{ We know that } |n\rangle\langle n| &= \frac{A^{+n}}{\sqrt{n!}} |0\rangle\langle 0| \frac{A^n}{\sqrt{n!}} \\ &= \frac{1}{n!} A^{+n} e^{-A^+; A} A^n \end{aligned}$$

(see the Exercise on page 112), so that

$$\begin{aligned} F(A^+A) &= \sum_{n=0}^{\infty} |n\rangle F(n) \langle n| \\ &= \sum_{n=0}^{\infty} \frac{F(n)}{n!} A^{+n} e^{-A^+; A} A^n, \end{aligned}$$

indeed.

$$\text{Now, for } F(A^+A) = z^{A^+A},$$

$$\begin{aligned} z^{A^+A} &= \sum_{n=0}^{\infty} \frac{1}{n!} z^n A^{+n} e^{-A^+; A} A^n \\ &= \left( e^{z A^+A} e^{-A^+A} \right)_{\text{normally ordered}} \\ &= e^{(z-1)A^+; A} \end{aligned}$$

 $\boxed{4}$  Writing

$$\begin{aligned} H &= \left( \frac{1}{2M} p_1^2 + \frac{1}{2} M \omega^2 x_1^2 \right) + \left( \frac{1}{2M} p_2^2 + \frac{1}{2} M \omega^2 x_2^2 \right) \\ &\quad + \left( \frac{1}{2M} p_3^2 + \frac{1}{2} M \omega^2 x_3^2 \right) \end{aligned}$$

We see that the joint eigenstates of the three individual one-dimensional harmonic oscillators are eigenstates of  $H$ , so that

$$H |n_1, n_2, n_3\rangle = \hbar \omega \left( n_1 + n_2 + n_3 + \frac{3}{2} \right),$$

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$$\text{or } H|n_1, n_2, n_3\rangle = |n_1, n_2, n_3\rangle h\omega(N + \frac{3}{2})$$

with  $N = n_1 + n_2 + n_3$ . To given  $N$  value we have

$$n_1 = 0, 1, 2, \dots, N;$$

$$n_2 = 0, 1, \dots, N - n_1;$$

$$n_3 = N - n_1 - n_2.$$

So for each choice of both  $n_1$  and  $n_2$ , there is just one value of  $n_3$ , one possible ket  $|n_1, n_2, n_3\rangle$ . And for a given value of  $n_1$ , we have  $N - n_1 + 1$  possible values of  $n_2$ . The total count of kets  $|n_1, n_2, n_3\rangle$  with given  $N = n_1 + n_2 + n_3$  is therefore

$$\sum_{n_1=0}^N (N - n_1 + 1) = \sum_{m=1}^{N+1} m = \frac{1}{2}(N+1)(N+2).$$

[5] Using  $L_1 = \frac{1}{2}(L_+ + L_-)$ ,  $L_2 = \frac{1}{2i}(L_+ - L_-)$ , we have

$$L_1|1\rangle = |l=1, m=1\rangle \frac{\hbar}{3\sqrt{2}} + |l=1, m=0\rangle \frac{\sqrt{8}\hbar}{3} + |l=1, m=-1\rangle \frac{\hbar}{3\sqrt{2}}$$

$$L_2|1\rangle = |l=1, m=1\rangle \frac{\hbar}{3\sqrt{2}} + |l=1, m=-1\rangle \frac{i\hbar}{3\sqrt{2}}$$

so that

$$\langle L_1 \rangle = \frac{2}{3} \frac{\hbar}{3\sqrt{2}} + \frac{1}{3} \frac{\sqrt{8}\hbar}{3} + \frac{2}{3} \frac{\hbar}{3\sqrt{2}} = \frac{4\sqrt{2}}{9} \hbar$$

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$$\text{and } \langle L_2 \rangle = \frac{2}{3} \frac{\hbar}{i3\sqrt{2}} + \frac{2}{3} \frac{i\hbar}{3\sqrt{2}} = 0$$

as well as

$$\langle L_1^2 \rangle = \left(\frac{\hbar}{3\sqrt{2}}\right)^2 + \left(\frac{\sqrt{2}\hbar}{3}\right)^2 + \left(\frac{\hbar}{3\sqrt{2}}\right)^2 = \hbar^2$$

$$\text{and } \langle L_2^2 \rangle = \left(\frac{\hbar}{3\sqrt{2}}\right)^2 + \left(\frac{\hbar}{3\sqrt{2}}\right)^2 = \frac{1}{9} \hbar^2,$$

$$\text{giving } \delta L_1 = \sqrt{\hbar^2 - \frac{3^2}{81} \hbar^2} = \frac{7}{9} \hbar$$

$$\text{and } \delta L_2 = \sqrt{\frac{1}{9} \hbar^2 - 0} = \frac{1}{3} \hbar.$$

For  $L_3$  we find immediately that

$$\langle L_3 \rangle = \left(\frac{2}{3}\right)^2 \hbar + \left(\frac{1}{3}\right)^2 0\hbar + \left(\frac{2}{3}\right)^2 (-\hbar) = 0,$$

$$\langle L_3^2 \rangle = \left(\frac{2}{3}\right)^2 \hbar^2 + \left(\frac{1}{3}\right)^2 (0\hbar)^2 + \left(\frac{2}{3}\right)^2 (-\hbar)^2 = \frac{8}{9} \hbar^2,$$

$$\text{giving } \delta L_3 = \frac{2\sqrt{2}}{3} \hbar.$$

The uncertainty relation for  $L_1$  and  $L_2$  is

$$\delta L_1 \delta L_2 \geq \frac{1}{2} |i \langle [L_1, L_2] \rangle| = \frac{\hbar}{2} |\langle L_3 \rangle|,$$

here:

$$\frac{7}{9} \hbar \frac{1}{3} \hbar \geq \frac{\hbar}{2} 0, \text{ OK.}$$

$$\text{Likewise } \delta L_2 \delta L_3 \geq \frac{\hbar}{2} |\langle L_1 \rangle|,$$

here:

$$\frac{1}{3} \hbar \frac{2\sqrt{2}}{3} \hbar \geq \frac{\hbar}{2} \frac{4\sqrt{2}}{9} \hbar, \text{ OK with } = \text{ sign.}$$

$$\text{Finally } \delta L_3 \delta L_1 \geq \frac{\hbar}{2} |\langle L_2 \rangle|,$$

$$\text{here: } \frac{2\sqrt{2}}{3} \hbar \frac{7}{9} \hbar \geq \frac{\hbar}{2} 0, \text{ OK.}$$