

Question 1/6.....

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[1] Since $0 = [A, BB^{-1}] = B[A, B^{-1}] + [A, B]B^{-1}$
we have

$$[A, B^{-1}] = -B^{-1}[A, B]B^{-1}.$$

[2] The oscillator will remain in the instantaneous ground state. In view of

$$H = \hbar\omega \left(A^\dagger - \frac{\Omega^*}{\omega} \right) \left(A - \frac{\Omega}{\omega} \right) - \frac{1}{\hbar} \frac{|\Omega|^2}{\omega}$$

the instantaneous ground state is the ground state of an oscillator with
 $\langle A \rangle = \Omega/\omega$, $\langle A^\dagger \rangle = \Omega^*/\omega$, so that

$$\rho = e^{-\left(A^\dagger - \Omega^*/\omega \right); \left(A - \Omega/\omega \right)}$$

with $\Omega = \Omega(t)$ applies at intermediate times and

$$\rho(T) = e^{-\left(A^\dagger - \Omega_0^*/\omega \right); \left(A - \Omega_0/\omega \right)}$$

is the statistical operator (= projector to the ground state) at time T .

(a) With $A|0\rangle = 0$ and $\langle 0|A^\dagger = 0$, we have

$$\langle 0|\rho(T)|0\rangle = e^{-|\Omega_0|^2/\omega^2}$$

for the probability of finding the oscillator in the ground state $|0\rangle$ to $\Omega=0$, when looking for it at time T .

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(b) We have $|1\rangle = A^\dagger |0\rangle$, $\langle 1| = \langle 0|A$,
so that

$$\text{prob(in } |1\rangle \text{ at } T) = \langle 0|A \rho(T) A^\dagger |0\rangle$$

$$\text{where } A \rho(T) A^\dagger = AA^\dagger \rho(T) + A[\rho(T), A^\dagger]$$

$$= AA^\dagger \rho(T) + A \frac{\partial}{\partial A} \rho(T)$$

$$= AA^\dagger \rho(T) - A (A^\dagger - \Omega_\infty^*/\omega) \rho(T)$$

$$= (\Omega_\infty^*/\omega) A \rho(T)$$

$$= (\Omega_\infty^*/\omega) (\rho(T)A + [A, \rho(T)])$$

$$= (\Omega_\infty^*/\omega) (\rho(T)A + \frac{\partial}{\partial A^\dagger} \rho(T))$$

$$= (\Omega_\infty^*/\omega) (\rho(T)A - \rho(T)(A - \Omega_\infty/\omega))$$

$$= |\Omega_\infty/\omega|^2 \rho(T),$$

$$\text{so that } \langle 0|A \rho(T) A^\dagger |0\rangle$$

$$= |\Omega_\infty/\omega|^2 \langle 0|\rho(T)|0\rangle$$

$$= |\Omega_\infty/\omega|^2 e^{-|\Omega_\infty/\omega|^2}$$

Note: There are quite a few different methods for getting this result, such as making use of the explicit Gaussian wave functions.

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$$\boxed{3} \text{ We have } V_2(\vec{R}) = e^{-i\vec{a}\cdot\vec{P}/\hbar} V_1(\vec{R}) e^{i\vec{a}\cdot\vec{P}/\hbar}$$

$$= W^\dagger V_1(\vec{R}) W$$

with the unitary W commuting with $H_0 = \frac{1}{2M} \vec{P}^2$ and therefore also with $G = (E - H_0 + i\epsilon)^{-1}$. In the equation for the transition operator to V_2 ,

$$T_2 = V_2 + V_2 G T_2,$$

we have then

$$T_2 = W^\dagger V_1 W + W^\dagger V_1 G W T_2$$

or

$$W T_2 W^\dagger = V_1 + V_1 G W T_2 W^\dagger,$$

so that $W T_2 W^\dagger$ obeys the equation for the transition operator to V_1 :

$$T_1 = V_1 + V_1 G T_1.$$

It follows that $W T_2 W^\dagger = T_1$, or

$$T_2 = W^\dagger T_1 W.$$

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14(a) In $T = V + VG + VG^2 + \dots$, every term has one V to the left and one V to the right, so that $T|1\rangle$ is a linear combination of $|1_1\rangle$ and $|1_2\rangle$, and $\langle 1|T$ is a superposition of $\langle 1_1|$ and $\langle 1_2|$. Therefore,

$$T = (|1_1\rangle, |1_2\rangle) \begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix} \begin{pmatrix} \langle 1_1| \\ \langle 1_2| \end{pmatrix}.$$

(b) Taking matrix elements of $T = V + VG$ we get

$$\begin{aligned} T_{11} &= \langle 1_1 | T | 1_1 \rangle = V_1 + V_1 \langle 1_1 | G | 1_1 \rangle \\ &= V_1 + V_1 \langle 1_1 | G (|1_1\rangle T_{11} + |1_2\rangle T_{21}) \\ &= V_1 + V_1 (G_{11} T_{11} + G_{12} T_{21}) \end{aligned}$$

with $G_{11} = \langle 1_1 | G | 1_1 \rangle$, $G_{12} = \langle 1_1 | G | 1_2 \rangle$.
Likewise

$$T_{21} = V_2 G_{21} T_{11} + V_2 G_{22} T_{21},$$

$$T_{12} = V_1 G_{11} T_{12} + V_1 G_{12} T_{22},$$

$$T_{22} = V_2 + V_2 G_{21} T_{12} + V_2 G_{22} T_{22}.$$

We solve for T_{11} , T_{12} , T_{21} , and T_{22} and get

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$$\begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix} = \left[\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \begin{pmatrix} V_1 & 0 \\ 0 & V_2 \end{pmatrix} \begin{pmatrix} G_{11} & G_{12} \\ G_{21} & G_{22} \end{pmatrix} \right]^{-1} \begin{pmatrix} V_1 & 0 \\ 0 & V_2 \end{pmatrix}$$

$$= \frac{1}{\text{DET}} \begin{pmatrix} V_1 - V_1 V_2 G_{22} & V_1 G_{12} V_2 \\ V_2 G_{21} V_1 & V_2 - V_2 V_1 G_{11} \end{pmatrix}$$

with $\text{DET} = (1 - V_1 G_{11})(1 - V_2 G_{22}) - V_1 G_{12} V_2 G_{21}$.

5] When composing $j=2$ and $j=1$ from $j=3/2$ and $j=1/2$, we have $|j=m=2\rangle = |3/2, 1/2\rangle$
 short for $j=3/2, m=3/2$ \uparrow
 short for $j=1/2, m=1/2$ \uparrow

so that $|j=2, m=1\rangle \propto J_- |3/2, 1/2\rangle$

$$= |1/2, 1/2\rangle \hbar \sqrt{\left(\frac{3}{2} + \frac{3}{2}\right)\left(\frac{3}{2} - \frac{3}{2} + 1\right)}$$

$$+ |3/2, -1/2\rangle \hbar \sqrt{\left(\frac{1}{2} + \frac{1}{2}\right)\left(\frac{1}{2} - \frac{1}{2} + 1\right)}$$

$$= |1/2, 1/2\rangle \hbar \sqrt{3} + |3/2, -1/2\rangle \hbar.$$

The state with $j=1, m=1$ is orthogonal to this, but a superposition of the same two kets. Properly normalized, we have

$$|j=1, m=1\rangle = (|3/2, -1/2\rangle \sqrt{3} - |1/2, 1/2\rangle) / 2.$$

Another application of J_- gives

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$$\begin{aligned} J_- |j=1, m=1\rangle &= |j=1, m=0\rangle \hbar \sqrt{(1+1)(1-1+1)} \\ &= (J_- | \frac{3}{2}; -\frac{1}{2} \rangle \sqrt{3} - J_- | \frac{1}{2}; \frac{1}{2} \rangle) / 2 \\ &= (| \frac{1}{2}; -\frac{1}{2} \rangle \hbar \sqrt{3}^2 - | -\frac{1}{2}; \frac{1}{2} \rangle \hbar \sqrt{4} \\ &\quad - | \frac{1}{2}; -\frac{1}{2} \rangle \hbar) / 2 \\ &= (| \frac{1}{2}; -\frac{1}{2} \rangle - | -\frac{1}{2}; \frac{1}{2} \rangle) \hbar, \end{aligned}$$

So that $|j=1, m=1\rangle = (| \frac{1}{2}; -\frac{1}{2} \rangle - | -\frac{1}{2}; \frac{1}{2} \rangle) / \sqrt{2}$

m_1 m_2

(a) These probabilities are

0 for $m_2 = \frac{3}{2}$ and $m_2 = -\frac{3}{2}$,

$\frac{1}{2}$ for $m_2 = \frac{1}{2}$ and $m_2 = -\frac{1}{2}$.

(b) These probabilities are $\frac{1}{2}$ each.

(c) This probability is $\frac{1}{2}$.