

Question... 1/5

Do not write on either margin

Write answers on this side of the paper only.

$$\square (a) \frac{d}{dt} X = \frac{\partial H}{\partial P} = v,$$

$$X(t) = X(t_0) + vT \text{ with } T = t - t_0,$$

$$\frac{d}{dt} P = -\frac{\partial H}{\partial X} = -k^2 X$$

$$P(t) = P(t_0) - k^2 T X(t_0) - \frac{1}{2} k^2 v T^2.$$

$$(b) i\hbar \frac{\partial}{\partial t} \langle x, t | p, t_0 \rangle = \langle x, t | H | p, t_0 \rangle$$

$$\text{with } H = vP(t) + \frac{1}{2} k^2 X(t)^2$$

$$= vP(t_0) + \frac{1}{2} k^2 X(t_0)^2$$

$$= vP(t_0) + \frac{1}{2} k^2 (X(t) - vT)^2,$$

So that

$$i\hbar \frac{\partial}{\partial t} \log \langle x, t | p, t_0 \rangle$$

$$= vP + \frac{1}{2} k^2 (X - vT)^2$$

$$= \frac{\partial}{\partial t} \left(vTP - \frac{k^2}{6v} (X - vT)^3 \right)$$

implying

$$\langle x, t | p, t_0 \rangle = \frac{e^{ixp/\hbar}}{\sqrt{2\pi\hbar}} e^{-ivTp/\hbar} e^{\frac{ik^2}{6\hbar v} [(X - vT)^3 - X^3]}$$

where the prefactor ensures the correct value for $t = t_0$.

Question. 2/5.....

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(c) Given the solutions of the equations of motion in (a), we have immediately

$$\langle X(t) \rangle = vT, \quad \langle P(t) \rangle = -\frac{1}{2}k^2vT^2,$$

and

$$\begin{aligned} \langle X(t)^2 \rangle &= \langle X(t_0)^2 \rangle + 2vT \langle X(t_0) \rangle + (vT)^2 \\ &= x_0^2 + (vT)^2, \end{aligned}$$

so that

$$\begin{aligned} \delta X(t) &= \sqrt{\langle X(t)^2 \rangle - \langle X(t) \rangle^2} \\ &= \sqrt{x_0^2} = |x_0| \stackrel{\text{assume } x_0 > 0}{=} x_0. \end{aligned}$$

Further,

$$\begin{aligned} \langle P(t)^2 \rangle &= \langle P(t_0)^2 \rangle + (k^2T)^2 \langle X(t_0)^2 \rangle \\ &\quad + \left(\frac{1}{2}k^2vT^2\right)^2 \\ &\quad - k^2T \langle (X(t_0)P(t_0) + P(t_0)X(t_0)) \rangle \\ &\quad - k^2vT^2 \langle P(t_0) \rangle \\ &\quad + k^4vT^3 \langle X(t_0) \rangle \\ &= p_0^2 + (k^2Tx_0)^2 + \left(\frac{1}{2}k^2vT^2\right)^2 \end{aligned}$$

giving

$$\delta P(t) = \sqrt{p_0^2 + (k^2Tx_0)^2}.$$

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[2] Multiply by coherent state bra $\langle a'^* |$:

$$\begin{aligned} \langle a'^* | e^{\zeta A^+} | a \rangle &= e^{\zeta a'^*} \langle a'^* | a \rangle \\ &= e^{\zeta a'^*} e^{a'^* a} = e^{a'^*(a+\zeta)} = \langle a'^* | a+\zeta \rangle \end{aligned}$$

and since the $\langle a'^* |$ bras are complete,
we get

$$e^{\zeta A^+} | a \rangle = | a+\zeta \rangle.$$

Similarly, $\langle a^* | e^{\zeta^* A} = \langle a^* + \zeta^* |$.

$$\begin{aligned} [3] \quad \vec{R} \times \vec{F} &= \vec{R} \times \frac{d}{dt} \vec{P} = \frac{d}{dt} (\vec{R} \times \vec{P}) - \frac{d\vec{R}}{dt} \times \vec{P} \\ &= \frac{d}{dt} (\vec{R} \times \vec{P}) - \frac{1}{M} \underbrace{\vec{P} \times \vec{P}}_{=0} = \frac{d}{dt} (\vec{R} \times \vec{P}) \\ &= \frac{1}{i\hbar} [\vec{R} \times \vec{P}, H] \end{aligned}$$

and since the expectation value
is taken for an eigenstate of
 H itself, we have

$$\langle \vec{R} \times \vec{F} \rangle = 0.$$

$$[4] (a) \text{ Since } L_1 \cos \gamma + L_2 \sin \gamma = \begin{pmatrix} L_1 \\ L_2 \\ L_3 \end{pmatrix} \cdot \underbrace{\begin{pmatrix} \cos \gamma \\ \sin \gamma \\ 0 \end{pmatrix}}_{\text{unit vector}},$$

it is a cartesian component of \vec{L}
in the l_2 -plane ($\equiv xy$ plane) and,

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therefore, the eigenvalues are the same as the eigenvalues of L_3 , namely $0, \pm \hbar$.

$$(b) \text{ We have } L_1 = \frac{1}{2}(L_+ + L_-), L_2 = \frac{1}{2i}(L_+ - L_-),$$

$$L_1 |1,1\rangle = |1,0\rangle \hbar/\sqrt{2},$$

$$L_2 |1,1\rangle = |1,0\rangle i\hbar/\sqrt{2},$$

$$L_1 |1,0\rangle = (|1,1\rangle + |1,-1\rangle) \hbar/\sqrt{2},$$

$$L_2 |1,0\rangle = (|1,1\rangle - |1,-1\rangle) (-i\hbar/\sqrt{2}),$$

$$L_1 |1,-1\rangle = |1,0\rangle \hbar/\sqrt{2},$$

$$L_2 |1,-1\rangle = |1,0\rangle (-i\hbar/\sqrt{2}), \text{ so that}$$

$$(L_1 \cos \gamma + L_2 \sin \gamma) |1,1\rangle = |1,0\rangle \frac{\hbar}{\sqrt{2}} e^{i\gamma},$$

$$(L_1 \cos \gamma + L_2 \sin \gamma) |1,0\rangle = |1,1\rangle \frac{\hbar}{\sqrt{2}} e^{-i\gamma} + |1,-1\rangle \frac{\hbar}{\sqrt{2}} e^{i\gamma},$$

$$(L_1 \cos \gamma + L_2 \sin \gamma) |1,-1\rangle = |1,0\rangle \frac{\hbar}{\sqrt{2}} e^{-i\gamma}.$$

It follows that the eigenket to eigenvalue 0 is

$$|e_{\nu=0}\rangle = (|1,1\rangle e^{-i\gamma} - |1,-1\rangle e^{i\gamma})/\sqrt{2},$$

and the eigenkets to eigenvalues $\pm \hbar$ are

$$|e_{\nu=\pm 1}\rangle = (|1,1\rangle e^{-i\gamma} \pm |1,0\rangle \sqrt{2} + |1,-1\rangle e^{i\gamma})/2.$$

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5 We have $H = H_0 + H_1$, with $H_0 =$ Hamilton operator of hydrogenic atoms as in the notes, and

$$H_1 = V_1(|\vec{r}|) \quad \text{with}$$

$$V_1(r) = V_{\text{sphere}}(r) - V_{\text{point}}(r) = \begin{cases} \frac{Ze^2}{r} - \frac{Ze^2}{2b^3}(3b^2 - r^2) & \text{for } r < b, \\ 0 & \text{for } r > b. \end{cases}$$

Therefore, the 1st-order correction to the ground-state energy is

$$\begin{aligned} \langle 100 | H_1 | 100 \rangle &= \int (d\vec{r}) V_1(r) |\langle \vec{r} | 100 \rangle|^2 \\ &= \int_0^b dr r^2 V_1(r) 4 \left(\frac{Z}{a_0}\right)^3 e^{-2Zr/a_0} \underbrace{\left(\frac{1}{4\pi} R_{10}(r)\right)^2}_{\approx 1 \text{ since } b \ll a_0} \end{aligned}$$

$$= 4 \left(\frac{Z}{a_0}\right)^3 Ze^2 \int_0^b dr \left(r - \frac{3}{2} \frac{r^2}{b} + \frac{1}{2} \frac{r^4}{b^3} \right)$$

$$= 4 \left(\frac{Z}{a_0}\right)^3 Ze^2 \left(\frac{1}{2} b^2 - \frac{1}{2} b^2 + \frac{1}{10} b^2 \right)$$

$$= \frac{Z^2 e^2}{2a_0} \frac{4}{5} \left(\frac{b}{a_0/Z}\right)^2, \quad \text{and the corrected}$$

ground-state energy is

$$E_0 = - \frac{Z^2 e^2}{2a_0} \left(1 - \frac{4}{5} \left(\frac{b}{a_0/Z}\right)^2 \right).$$