

Question... 1/6

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□(a) We have $UV|v_e\rangle = |v_{e+1}\rangle v_e$ and therefore

$$\begin{aligned} (UV)^N |v_e\rangle &= |v_{e+N}\rangle v_e v_{e+1} \cdots v_{e+N-1} \\ &= |v_e\rangle e^{i \frac{2\pi}{N} (1 + (1+1) + \cdots + (1+N-1))} \\ &= |v_e\rangle e^{i \frac{2\pi}{N} \frac{N}{2} (2N-1)} \\ &= |v_e\rangle e^{i\pi(N-1)}, \end{aligned}$$

So that

$$(UV)^N = (-1)^{N-1}.$$

(b) Similarly, $U^m V^n |v_e\rangle = |v_{e+m}\rangle v_e^n$ implies

$$\begin{aligned} (U^m V^n)^N |v_e\rangle &= |v_{e+mN}\rangle v_e^n v_{e+m}^n \cdots v_{e+m(N-1)}^n \\ &= |v_e\rangle e^{i \frac{2\pi}{N} n \frac{N}{2} (2N-1)} \\ &= |v_e\rangle e^{i\pi mn(N-1)} \end{aligned}$$

and

$$(U^m V^n)^N = (-1)^{(N-1)mn}$$

follows.

(c) From $S|v_k\rangle = |u_k\rangle$ we get

$$S = \sum_R |u_R\rangle \langle v_R|.$$

And then

$$\langle u_R | S = \langle v_R |,$$

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as well as

$$\begin{aligned}
 S|u_k\rangle &= \sum_e |u_e\rangle \langle v_e | u_k \rangle \\
 &= \sum_e |u_e\rangle \frac{1}{\sqrt{N}} e^{-i\frac{2\pi}{N}ek} \\
 &= \sum_e |u_e\rangle \frac{1}{\sqrt{N}} e^{i\frac{2\pi}{N}e(N-k)} \\
 &= \sum_e |u_e\rangle \langle u_e | v_{N-k} \rangle,
 \end{aligned}$$

so that

$$S|u_k\rangle = |v_{N-k}\rangle,$$

and likewise $\langle v_k | S = \langle u_{N-k} |$.

$$\begin{aligned}
 (d) \quad US &= U \sum_k |u_k\rangle \langle v_k | \\
 &= \sum_k |u_k\rangle u_k \langle v_k | \\
 &= \sum_k |u_k\rangle v_k \langle v_k | \\
 &= \sum_k |u_k\rangle \langle v_k | V = SV.
 \end{aligned}$$

$$\begin{aligned}
 (e) \quad \text{It is } S|v_k\rangle &= |u_k\rangle, \quad S|u_k\rangle = |v_{N-k}\rangle, \\
 S|v_{N-k}\rangle &= |u_{N-k}\rangle, \quad S|u_{N-k}\rangle = |v_k\rangle,
 \end{aligned}$$

so that $S^4|v_k\rangle = |v_k\rangle$ for all k . This states that $S^4 = 1$, so that S has period 4, except for $N=2$, where $S^2|v_1\rangle = |v_{2-1}\rangle = |v_1\rangle$, $S^2|v_2\rangle = |v_2\rangle$, and the period is 2.

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2 (a) For $v=0$, we have

$$i\hbar \frac{\partial}{\partial t_1} \langle x, t_1 | p, t_2 \rangle = \langle x, t_1 | (-Fx(t_1)) | p, t_2 \rangle$$

$$= -Fx \langle x, t_1 | p, t_2 \rangle$$

so that

$$\langle x, t_1 | p, t_2 \rangle = \langle x, t_2 | p, t_2 \rangle e^{\frac{i}{\hbar} Fx(t_1 - t_2)}$$

or

$$\langle x, t_1 | p, t_2 \rangle = \frac{e^{i x p / \hbar}}{\sqrt{2\pi\hbar}} e^{\frac{i}{\hbar} FxT}$$

with $T = t_1 - t_2$.

(b) Since $\frac{\partial}{\partial v} H = |P(t)|$, we have

$$\frac{\partial}{\partial v} \langle x, t_1 | p, t_2 \rangle = -\frac{i}{\hbar} \langle x, t_1 | \int_{t_2}^{t_1} dt |P(t)| | p, t_2 \rangle$$

where $P(t)$ obeys

$$\frac{d}{dt} P = -\frac{\partial H}{\partial x} = F,$$

so that

$$P(t) = P(t_2) + F(t - t_2),$$

implying

$$\frac{\partial}{\partial v} \log \langle x, t_1 | p, t_2 \rangle = -\frac{i}{\hbar} \int_{t_2}^{t_1} dt |p + F(t - t_2)|$$

$$= -\frac{i}{\hbar F} \int_p^{p+FT} dy |y| = -\frac{i}{2\hbar F} ((p+FT)|p+FT| - p|p|).$$

$p + F(t - t_2) = y$

It follows that

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$$\langle x, t_1 | p, t_2 \rangle = \langle x, t_1 | p, t_2 \rangle \Big|_{v=0} e^{-\frac{i v}{2 \hbar F} (|p+FT| |p+FT| - |p| |p|)}$$

$$= \frac{e^{i x p / \hbar}}{\sqrt{2 \pi \hbar}} e^{-\frac{i v}{2 \hbar F} [(p+FT) |p+FT| - |p| |p|]} + \frac{i}{\hbar} F x T$$

3(a) $\int_0^\pi \frac{d\vartheta}{\sin \vartheta} \langle \lambda | \vartheta \rangle \langle \vartheta | \lambda' \rangle$

$$= \int_0^\pi \frac{d\vartheta}{\sin \vartheta} \frac{1}{2\pi} \left(\tan \frac{\vartheta}{2} \right)^{-i(\lambda - \lambda')}$$

substitute $e^y = \tan \frac{\vartheta}{2}$, $\vartheta=0: y = -\infty$
 $\vartheta=\pi: y = \infty$

$$dy e^y = \frac{d\vartheta}{2(\cos \frac{\vartheta}{2})^2}, \quad dy = \frac{d\vartheta}{2 \tan \frac{\vartheta}{2} (\cos \frac{\vartheta}{2})^2}$$

$$= \frac{d\vartheta}{\sin \vartheta}$$

to get

$$2\pi \int_0^\pi \frac{d\vartheta}{\sin \vartheta} \langle \lambda | \vartheta \rangle \langle \vartheta | \lambda' \rangle = \int_{-\infty}^{\infty} dy e^{-iy(\lambda - \lambda')}$$

$$= 2\pi \delta(\lambda - \lambda') = 2\pi \langle \lambda | \lambda' \rangle,$$

and

$$\int_0^\pi \frac{d\vartheta}{\sin \vartheta} |\vartheta\rangle \langle \vartheta| = 1 \text{ follows.}$$

Similarly,

$$\int_{-\infty}^{\infty} d\lambda \langle \vartheta | \lambda \rangle \langle \lambda | \vartheta' \rangle$$

$$= \int_{-\infty}^{\infty} d\lambda \frac{1}{2\pi} \left(\frac{\tan \frac{\vartheta}{2}}{\tan \frac{\vartheta'}{2}} \right)^{i\lambda}$$

Question... 5/6

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$$= \delta \left(\log \frac{\tan \frac{\theta}{2}}{\tan \frac{\theta'}{2}} \right)$$

$$= \delta(\theta - \theta') \left| \left(\frac{\partial}{\partial \theta} \frac{\tan \frac{\theta}{2}}{\tan \frac{\theta'}{2}} \right) \Big|_{\theta = \theta'} \right|^{-1}$$

$$= \delta(\theta - \theta') \left| \frac{1}{2(\cos \frac{\theta}{2})^2 \tan \frac{\theta'}{2}} \right|_{\theta = \theta'}^{-1}$$

$$= \delta(\theta - \theta') \sin \theta = \langle \theta | \theta' \rangle,$$

implying $\int_{-\infty}^{\infty} d\lambda |\lambda\rangle \langle \lambda| = 1.$

(b) Since $f(\theta) = \int_0^{\pi} \frac{d\theta'}{\sin \theta'} |\theta'\rangle f(\theta') \langle \theta'|$ for any function of θ , we have

$$\begin{aligned} \left(\tan \frac{\theta}{2} \right)^{i\lambda'} |\lambda'\rangle &= \int_0^{\pi} \frac{d\theta'}{\sin \theta'} |\theta'\rangle \underbrace{\left(\tan \frac{\theta}{2} \right)^{i(\lambda'+\theta')}}_{\langle \theta' | \lambda'+\theta' \rangle} \\ &= |\lambda'+\theta'\rangle. \end{aligned}$$

(c) We have $e^{i\mu\Lambda} \left(\tan \frac{\theta}{2} \right)^{i\lambda} |\lambda'\rangle$

$$= e^{i\mu\Lambda} |\lambda+\lambda'\rangle = |\lambda+\lambda'\rangle e^{i\mu(\lambda+\lambda')}$$

$$= \left(\tan \frac{\theta}{2} \right)^{i\lambda} |\lambda'\rangle e^{i\mu(\lambda+\lambda')}$$

$$= \left(\tan \frac{\theta}{2} \right)^{i\lambda} e^{i\mu\Lambda} |\lambda'\rangle e^{i\mu\lambda}$$

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which implies

$$e^{i\mu\Lambda} \left(\tan \frac{\theta}{2}\right)^{i\lambda} = e^{i\mu\lambda} \left(\tan \frac{\theta}{2}\right)^{i\lambda} e^{i\mu\Lambda}$$

(d) Here we have

$$\begin{aligned} \langle \mathcal{J}' | &= \langle \mathcal{J} | e^{i\mu\Lambda} \\ &= e^{i\mu\lambda} \langle \mathcal{J} | \left(\tan \frac{\theta}{2}\right)^{i\lambda} e^{i\mu\Lambda} \left(\tan \frac{\theta}{2}\right)^{-i\lambda} \\ &= e^{i\mu\lambda} \left(\tan \frac{\theta}{2}\right)^{i\lambda} \left(\tan \frac{\theta'}{2}\right)^{-i\lambda} \langle \mathcal{J}' | \\ &= \underbrace{\left(e^{\mu} \left(\tan \frac{\theta}{2}\right) \left(\tan \frac{\theta'}{2}\right)^{-1}\right)^{i\lambda}}_{=1 \text{ for all } \lambda} \langle \mathcal{J}' | \end{aligned}$$

so that

$$\tan \frac{\theta'}{2} = e^{\mu} \tan \frac{\theta}{2}$$

or

$$\theta' = 2 \arctan \left(e^{\mu} \tan \frac{\theta}{2} \right).$$