### NATIONAL UNIVERSITY OF SINGAPORE

## PC2130 - Quantum Mechanics I

(Semester I: AY 13-14)

Examiner: Dr. Vitor M. Pereira

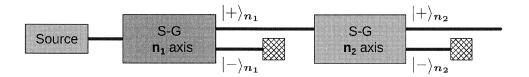
Time Allowed: 2 Hours

#### Instructions to candidates

- a) Please write your matriculation/registration number only. Do not write your name.
- b) This examination paper contains FOUR (4) problems and comprises FOUR (4) printed pages.
- c) This is a **closed book** examination.
- d) Answer **ALL** questions.
- e) Write your responses in the answer books only. Please start each problem in a new page.
- f) Percentage points are indicated beside the header of each problem, together with the partial points for each sub-question.
- g) No calculators or any other electronic device is allowed during the exam.
- h) Students are allowed to bring in an A4-sized (1 side only) sheet of manuscript notes.
- i) All answers should be adequately justified.

## **Problem 1** 25 % [5, 5, 5, 10]

Consider a beam of neutral atoms with total spin S=1/2 which is made to go through the sequence of Stern-Gerlach (SG) selectors depicted in the figure below. SG-n represents a Stern-Gerlach device oriented with its field gradient parallel to the direction of the unit vector n that selects the positive eigenstate,  $|+\rangle_n$ , of  $\hat{\mathbf{S}} \cdot \mathbf{n}$ , where  $\hat{\mathbf{S}}$  is the spin operator.



a) Expressing a unit vector n in terms of its spherical angles as

$$n = \sin \theta \cos \varphi \, u_x + \sin \theta \sin \varphi \, u_y + \cos \theta \, u_z,$$

write down a normalized state,  $|+\rangle_n$ , expressed in the eigenbasis of  $\hat{S}_z$  that represents the state of the atoms when they leave a SG-n selector.

b) For a general state of a quantum system,  $|\psi\rangle$ , one can define the density operator

$$\hat{\rho}_{\psi} = |\psi\rangle\langle\psi|$$
.

Show that the expectation value of any observable  $\hat{O}$  in the state  $|\psi\rangle$  can be written in terms of  $\hat{\rho}_{\psi}$  as

$$\langle \psi | \hat{O} | \psi \rangle = \text{Tr} \left( \hat{\rho}_{\psi} \hat{O} \right).$$

c) Show that the density operator,  $\hat{\rho}_n = |+\rangle_{nn} \langle +|$ , associated with the state of the atoms upon leaving a selector SG-n can be written as

$$\hat{\rho}_{\boldsymbol{n}} = \frac{1}{2} + \frac{1}{\hbar} \hat{\boldsymbol{S}} \cdot \boldsymbol{n}.$$

Suggestion: consider the action of  $\hat{\rho}_n$  on an arbitrary state expanded in the eigenbasis of  $\hat{S} \cdot n$ .

d) Consider now the case depicted in the figure above. Using the results above, or otherwise, show that, of all the atoms entering the second selector oriented along  $n_2$ , the fraction p leaving it can be written as

$$p=\frac{1+\boldsymbol{n}_1\cdot\boldsymbol{n}_2}{2},$$

where  $n_1$  and  $n_2$  are unit vectors defining the orientation of the two selectors.

Hint: it might be useful to recall that Pauli matrices are traceless and, for any two vectors  $\mathbf{a}$  and  $\mathbf{b}$ , the following identity involving the vector of Pauli matrices,  $\boldsymbol{\sigma}$ , holds:  $(\boldsymbol{\sigma} \cdot \mathbf{a}) (\boldsymbol{\sigma} \cdot \mathbf{b}) = (\mathbf{a} \cdot \mathbf{b}) \mathbf{1} + i (\mathbf{a} \times \mathbf{b}) \cdot \boldsymbol{\sigma}$ .

## **Problem 2** 20 % [20]

The time-evolved state of a quantum mechanical system can be generically written in terms of the evolution operator,  $\hat{U}(t,0)$ , as

$$|\psi(t)\rangle = \hat{U}(t,0) |\psi(0)\rangle$$
.

In the position representation this equation can be written as

$$\psi(x,t) = \int_{-\infty}^{+\infty} K(x,t; x', 0) \, \psi(x', 0) \, dx',$$

allowing us to calculate the wavefunction  $\psi(x,t)$  from the initial one,  $\psi(x,0)$ . Calculate the function K(x,t;x',0) for a free particle in 1D.

Note: the following result might be useful:

$$\int_{-\infty}^{+\infty} e^{-iax^2 + ibx} dx = \sqrt{\frac{\pi}{ia}} e^{\frac{ib^2}{4a}}, \qquad (a > 0).$$

# **Problem 3** 30 % [15, 15]

Consider a particle of mass m confined by the following infinite potential well:

$$V(x) = \begin{cases} +\infty, & x < -a \\ 0, & -a < x < a \\ +\infty, & x > a \end{cases}$$

Suddenly, this potential well is suppressed and immediately replaced by a new one:

$$V(x) = -aV_0 \,\delta(x)$$
  $(a, V_0 > 0).$ 

If the energy is measured right after this new potential is turned on, what is the probability of obtaining the value  $E=-\frac{ma^2V_0^2}{2\hbar^2}$  when:

- a) the particle was initially in the ground state of the infinite potential well?
- b) the particle was initially in the first excited state of the infinite potential well?

Note: you do not need to explicitly evaluate integrals, and can express your final result in terms of  $I_a(\alpha, \beta)$  defined as

$$I_a(\alpha, \beta) = \int_0^a \cos(\alpha x) \exp(-\beta x) dx.$$

**Problem 4** 25 % [5, 5, 5, 5, 2+3]

Consider the Hamiltonian for the 1D harmonic oscillator,

$$\hat{H} = \frac{\hat{P}^2}{2m} + \frac{1}{2}m\omega^2 \hat{X}^2,$$

and denote its eigenstates by  $|\varphi_n\rangle$ . The annihilation (or destruction) and number operators are defined as

$$a = \sqrt{\frac{m\omega}{2\hbar}}\,\hat{X} + i\sqrt{\frac{1}{2m\hbar\omega}}\,\hat{P}, \qquad \hat{N} = a^{\dagger}a.$$

- a) Use the definition of a above and its action on the eigenstates  $|\varphi_n\rangle$  to derive the normalized ground state wavefunction,  $\varphi_0(x) = \langle x|\varphi_0\rangle$ .
- b) Calculate the product of uncertainties  $\delta X_n \delta P_n$  when the system is in one of its stationary states  $|\varphi_n\rangle$ , expressing it as a function of the quantum number n.
- c) In view of Heisenberg's uncertainty principle, explain why the *exact* value of the product  $\delta X_0 \delta P_0$  in the ground state of the harmonic oscillator can be anticipated without its explicit calculation.
- d) Calculate the expectation value of the position,  $\langle \hat{X} \rangle$ , as a function of time when the initial state of the system is

$$|\psi(t=0)\rangle = \frac{1}{\sqrt{2}} |\varphi_0\rangle + \frac{1}{\sqrt{2}} |\varphi_1\rangle.$$

e) Consider the following identity regarding a unitary transformation (a translation) of the position operator

$$e^{i\hat{P}s/\hbar}\hat{X}e^{-i\hat{P}s/\hbar} = \hat{X} + s \qquad (s \in \mathbb{R}).$$

- i) From this identity, or otherwise, calculate  $\left[\hat{X},e^{i\hat{P}s/\hbar}\right]$  and  $\left[a,e^{i\hat{P}s/\hbar}\right]$ .
- ii) Use the previous commutator to show that

$$\langle \varphi_n | e^{i\hat{P}s/\hbar} | \varphi_0 \rangle = C_n \langle \varphi_0 | e^{i\hat{P}s/\hbar} | \varphi_0 \rangle,$$

and write the value of the constant  $C_n$ .

Note: the following result might be useful for this problem:

$$\int_{-\infty}^{+\infty} e^{-ax^2} dx = \sqrt{\frac{\pi}{a}}, \qquad (\operatorname{Re} a > 0).$$