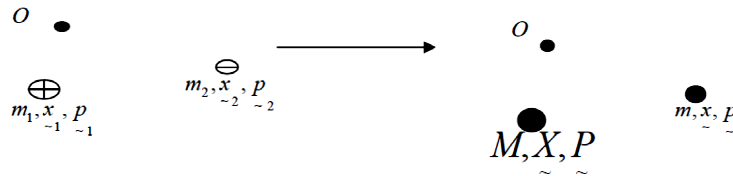


Question 1(a)

The Hamiltonian for the hydrogen atom, one proton and one electron with masses m_p and m_e ,

$$H = \frac{p_p^2}{2m_p} + \frac{p_e^2}{2m_e} + V(\vec{x}_p - \vec{x}_e)$$

We transform them into 2 equivalent particles M and m :



We let

$$M = m_p + m_e \approx m_p, \quad m = \frac{m_p m_e}{m_p + m_e} \approx m_e$$

Then we get

$$M\vec{X} = m_p\vec{x}_p + m_e\vec{x}_e, \quad \vec{x} = \vec{x}_p - \vec{x}_e$$

The momentum,

$$\vec{P} = M\dot{\vec{X}}, \quad \vec{p} = m\dot{\vec{x}}$$

Rearranging, we get

$$\vec{p}_p = \frac{m}{m_e}\vec{P} + \vec{p} \approx \vec{P} + \vec{p}, \quad \vec{p}_e = \frac{m}{m_p}\vec{P} - \vec{p} \approx -\vec{p}$$

Using the above relations, the Hamiltonian now becomes

$$\begin{aligned} H &= \frac{1}{2m_p} \left(\frac{m}{m_e} \vec{P} + \vec{p} \right)^2 + \frac{1}{2m_e} \left(\frac{m}{m_p} \vec{P} - \vec{p} \right)^2 + V(\vec{x}) \\ &= \frac{1}{2m_p} \left[\frac{m^2}{m_e^2} P^2 + \frac{m}{m_e} (\vec{P} \cdot \vec{p} + \vec{p} \cdot \vec{P}) + p^2 \right] + \frac{1}{2m_e} \left[\frac{m^2}{m_p^2} P^2 + \frac{m}{m_e} (\vec{P} \cdot \vec{p} - \vec{p} \cdot \vec{P}) + p^2 \right] + V(\vec{x}) \\ &= \frac{P^2}{2(m_p + m_e)} + \frac{p^2}{2} \left(\frac{m_p + m_e}{m_p m_e} \right) + V(\vec{x}) \\ &= \frac{P^2}{2M} + \frac{p^2}{2m} + V(\vec{x}) \end{aligned}$$

Where M, P describes the motion of the center of mass, and m, p describes the relative motion of the electron and the proton. [shown]

Question 1(b)

We know that

$$\frac{\partial r}{\partial x_i} = \frac{x_i}{r}, \quad \frac{\partial \theta}{\partial x_i} = \frac{1}{r \sin \theta} (x_i \cos \theta - \delta_{i3}), \quad \frac{\partial \phi}{\partial x_i} = \frac{1}{r \sin \theta} (\delta_{i2} \cos \phi - \delta_{i1} \sin \phi)$$

Therefore we have

$$\begin{aligned} l_3 &= x_1 p_2 - x_2 p_1 \\ &= \frac{\hbar}{i} \left(x_1 \frac{\partial}{\partial x_2} - x_2 \frac{\partial}{\partial x_1} \right) \\ &= \frac{\hbar}{i} \left[x_1 \left(\frac{\partial r}{\partial x_2} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x_2} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial x_2} \frac{\partial}{\partial \phi} \right) - x_2 \left(\frac{\partial r}{\partial x_1} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x_1} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial x_1} \frac{\partial}{\partial \phi} \right) \right] \end{aligned}$$

$$\begin{aligned}
&= \frac{\hbar}{i} \left[x_1 \left(\frac{x_2}{r} \frac{\partial}{\partial r} + \frac{x_2 \cos \theta}{r^2 \sin \theta} \frac{\partial}{\partial \theta} + \frac{\cos \phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) - x_2 \left(\frac{x_1}{r} \frac{\partial}{\partial r} + \frac{x_1 \cos \theta}{r^2 \sin \theta} \frac{\partial}{\partial \theta} - \frac{\sin \phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \right] \\
&= \frac{\hbar}{i} \left(\frac{x_1 \cos \phi}{r \sin \theta} \frac{\partial}{\partial \phi} + \frac{x_2 \cos \phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \\
&= \frac{\hbar}{i} \left(\cos^2 \phi \frac{\partial}{\partial \phi} + \sin^2 \phi \frac{\partial}{\partial \phi} \right) \\
&= \frac{\hbar}{i} \frac{\partial}{\partial \phi}
\end{aligned}$$

Let ψ be an arbitrary function, hence we have

$$\left[\phi, -i\hbar \frac{\partial}{\partial \phi} \right] \psi = -i\hbar \phi \frac{\partial \psi}{\partial \phi} + i\hbar \frac{\partial}{\partial \phi} (\phi \psi) = -i\hbar \phi \frac{\partial \psi}{\partial \phi} + i\hbar \psi + i\hbar \phi \frac{\partial \psi}{\partial \phi} = i\hbar \psi$$

$$\therefore \left[\phi, -i\hbar \frac{\partial}{\partial \phi} \right] = i\hbar \quad [\text{shown}]$$

Suppose $[\phi, l_3] = i\hbar$, then $\langle l, m | [\phi, l_3] | l, m \rangle = i\hbar \langle l, m | l, m \rangle$

$$LHS = \langle l, m | [\phi, l_3] | l, m \rangle = \langle l, m | (\phi l_3 - l_3 \phi) | l, m \rangle = \langle l, m | (\phi m \hbar - m \hbar \phi) | l, m \rangle = 0$$

$$RHS = i\hbar \langle l, m | l, m \rangle \neq 0$$

$$\therefore [\phi, l_3] \neq i\hbar$$

In coordinate representation,

$$l_+ = \hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right), \quad l_3 = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

With $l_3 Y_l^l(\theta, \phi) = l\hbar Y_l^l(\theta, \phi)$, we have

$$\frac{\hbar}{i} \frac{\partial}{\partial \phi} Y_l^l(\theta, \phi) = l\hbar Y_l^l(\theta, \phi), \quad \Rightarrow \quad Y_l^l(\theta, \phi) = f(\theta) e^{il\phi}$$

Where $f(\theta)$ is an arbitrary function of θ . Substitute into the representation of l_+ ,

$$\begin{aligned}
l_+ Y_l^l(\theta, \phi) &= 0 \\
\hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) f(\theta) e^{il\phi} &= 0 \\
\left(\frac{\partial}{\partial \theta} - l \cot \theta \right) f(\theta) &= 0 \\
\frac{\partial f}{\partial \theta} &= l f \cot \theta \\
\int \frac{1}{f} df &= \int l \cot \theta d\theta \\
\ln f &= l \ln(\sin \theta) + c \\
f &= N \sin^l \theta
\end{aligned}$$

$$\therefore Y_l^l(\theta, \phi) = N \sin^l \theta e^{il\phi} \quad [\text{shown}]$$

Question 2(a)

$$H = \frac{mv^2}{2} + V = \frac{(\vec{p} - q\vec{A})^2}{2m} + qA_0$$

Using the formula given, we have

$$\begin{aligned} \frac{d}{dt}\langle x_i \rangle &= \frac{i}{\hbar} \langle [H, x_i] \rangle + \underbrace{\langle \frac{\partial x_i}{\partial t} \rangle}_{=0} \\ &= \frac{i}{\hbar} \left\langle \left[\frac{(\vec{p} - q\vec{A})^2}{2m} + qA_0, x_i \right] \right\rangle \\ &= \frac{i}{2m\hbar} \langle [(p - qA)_j (p - qA)_j, x_i] \rangle \\ &= \frac{i}{2m\hbar} \langle (p - qA)_j [(p - qA)_j, x_i] + [(p - qA)_j, x_i] (p - qA)_j \rangle \\ &= \frac{i}{2m\hbar} \langle (p - qA)_j (-i\hbar\delta_{ij}) + (-i\hbar\delta_{ij})(p - qA)_j \rangle \\ &= \frac{i}{2m\hbar} \langle -2i\hbar(p - qA)_j \rangle \\ &= \frac{1}{m} \langle (p - qA)_j \rangle \end{aligned}$$

$$\therefore \frac{d}{dt}\langle \vec{x} \rangle = \frac{1}{m} \langle (\vec{p} - q\vec{A}) \rangle$$

Question 2(b)(i)

We know that $\vec{B} = \vec{\nabla} \times \vec{A}$. Then also,

$$\vec{A} = \frac{B_0}{2} (x_1\hat{j} - x_2\hat{i}) = \frac{1}{2} \vec{B} \times \vec{x}, \quad \vec{x} = x_1\hat{i} + x_2\hat{j} + x_3\hat{k}, \quad \vec{B} = B_0\hat{k}$$

We substitute \vec{A} into the formula,

$$\begin{aligned} \vec{\nabla} \times \vec{A} &= \vec{\nabla} \times \left(\frac{1}{2} \vec{B} \times \vec{x} \right) \\ &= \frac{1}{2} [(\vec{x} \cdot \vec{\nabla})\vec{B} - (\vec{B} \cdot \vec{\nabla})\vec{x} + \vec{B}(\vec{\nabla} \cdot \vec{x}) - \vec{x}(\vec{\nabla} \cdot \vec{B})] \\ &= \frac{1}{2} [0 - B_0\hat{k} + 3B_0\hat{k} - 0] \\ &= B_0\hat{k} \\ &= \vec{B} \end{aligned}$$

Next, we know that $\vec{E} = -2Cx_3\hat{k}$. Substituting A_0 into the formula

$$-\vec{\nabla}A_0 + \underbrace{\frac{\partial \vec{A}}{\partial t}}_{=0} = -\vec{\nabla}(Cx_3^2) = -2Cx_3\hat{k}$$

$$\therefore \vec{A} = \frac{B_0}{2} (x_1\hat{j} - x_2\hat{i}), \quad A_0 = Cx_3^2 \quad [\text{shown}]$$

Question 2(b)(ii)

The Hamiltonian now becomes

$$\begin{aligned}
 H &= \frac{(\vec{p} - q\vec{A})^2}{2m} + qA_0 \\
 &= \frac{\left[\vec{p} - q\frac{B_0}{2}(x_1\hat{j} - x_2\hat{i})\right]^2}{2m} + qCx_3^2 \\
 &= \frac{1}{2m} \left[p^2 + \frac{q^2 B_0^2}{4} (x_1^2 + x_2^2) - qB_0 l_3 \right] + qCx_3^2 \\
 &= \left[\frac{1}{2m} (p_1^2 + p_2^2) + \frac{q^2 B_0^2}{8m} (x_1^2 + x_2^2) \right] + \left[\frac{p_3^2}{2m} + qCx_3^2 \right] - \frac{qB_0 l_3}{2m} \\
 &= \left[\frac{1}{2m} (p_1^2 + p_2^2) + \frac{1}{2} m \underbrace{\left(\frac{q^2 B_0^2}{4m^2} \right)}_{=\omega_1^2} (x_1^2 + x_2^2) \right] + \left[\frac{p_3^2}{2m} + \frac{1}{2} m \underbrace{\left(\frac{2qC}{m} \right)}_{=\omega_2^2} x_3^2 \right] - \frac{qB_0 l_3}{2m}
 \end{aligned}$$

Which is a combination of 2 harmonic oscillators. We then find the energies to be

$$E(n_1, n_2) = \left(n_1 + \frac{1}{2}\right) \hbar\omega_1 + \left(n_2 + \frac{1}{2}\right) \hbar\omega_2 \quad [\text{shown}]$$

When $\vec{E} = 0$, we get the energy $E = \left(n_1 + \frac{1}{2}\right) \hbar\omega_1$, which is the Landau level of the particle.

Question 3(a)

$$\mathcal{R}_{\vec{u}}(\varepsilon) = 1 + \varepsilon\vec{u} \times + \frac{1}{2}\varepsilon^2\vec{u} \times \vec{u} \times + O(\varepsilon^3)$$

$$\mathcal{R}_{\vec{v}}\mathcal{R}_{\vec{u}} = \left(1 + \varepsilon\vec{v} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)\right) \left(1 + \varepsilon\vec{u} \times + \frac{1}{2}\varepsilon^2\vec{u} \times \vec{u} \times + O(\varepsilon^3)\right)$$

$$= 1 + \varepsilon(\vec{u} + \vec{v}) \times + \varepsilon^2\vec{v} \times \vec{u} \times + \frac{1}{2}\varepsilon^2\vec{u} \times \vec{u} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)$$

$$\mathcal{R}_{\vec{u}}^{-1}\mathcal{R}_{\vec{v}}\mathcal{R}_{\vec{u}} = \left(1 - \varepsilon\vec{u} \times + \frac{1}{2}\varepsilon^2\vec{u} \times \vec{u} \times + O(\varepsilon^3)\right) \left[1 + \varepsilon(\vec{u} + \vec{v}) \times + \varepsilon^2\vec{v} \times \vec{u} \times + \frac{1}{2}\varepsilon^2\vec{u} \times \vec{u} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)\right]$$

$$= 1 + \varepsilon\vec{v} \times + \varepsilon^2\vec{v} \times \vec{u} \times - \varepsilon^2\vec{u} \times \vec{v} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)$$

$$\mathcal{R}_{\vec{v}}^{-1}\mathcal{R}_{\vec{u}}^{-1}\mathcal{R}_{\vec{v}}\mathcal{R}_{\vec{u}} = \left(1 - \varepsilon\vec{v} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)\right) \left(1 + \varepsilon\vec{v} \times + \varepsilon^2\vec{v} \times \vec{u} \times - \varepsilon^2\vec{u} \times \vec{v} \times + \frac{1}{2}\varepsilon^2\vec{v} \times \vec{v} \times + O(\varepsilon^3)\right)$$

$$= 1 + \varepsilon^2\vec{v} \times \vec{u} \times - \varepsilon^2\vec{u} \times \vec{v} \times + O(\varepsilon^3)$$

$$\mathcal{R}_{\vec{v}}^{-1}\mathcal{R}_{\vec{u}}^{-1}\mathcal{R}_{\vec{v}}\mathcal{R}_{\vec{u}}\vec{x} = \vec{x} + \varepsilon^2\vec{v} \times (\vec{u} \times \vec{x}) - \varepsilon^2\vec{u} \times (\vec{v} \times \vec{x}) + O(\varepsilon^3)$$

$$= \vec{x} + \varepsilon^2[\vec{u}(\vec{v} \cdot \vec{x}) - \vec{v}(\vec{u} \cdot \vec{x})] + O(\varepsilon^3)$$

$$= \vec{x} + \varepsilon^2\vec{x} \times (\vec{u} \times \vec{v}) + O(\varepsilon^3)$$

$$= \vec{x} - \varepsilon^2\vec{w} \times \vec{x} + O(\varepsilon^3)$$

But we know that $\mathcal{R}_{\vec{w}}(-\varepsilon^2) = \vec{x} - \varepsilon^2\vec{w} \times \vec{x} + O(\varepsilon^4)$. $\therefore \mathcal{R}_{\vec{v}}^{-1}\mathcal{R}_{\vec{u}}^{-1}\mathcal{R}_{\vec{v}}\mathcal{R}_{\vec{u}}$ differs from $\mathcal{R}_{\vec{w}}(-\varepsilon^2)$ only by terms of higher order than ε^2 . [shown]

Question 3(b)

$$J_1^2 |j_1, m_1\rangle = \frac{3}{4} \hbar^2 |j_1, m_1\rangle = j_1(j_1 + 1) \hbar^2 |j_1, m_1\rangle$$

$$j_1^2 + j_1 - \frac{3}{4} = 0 \Rightarrow j_1 = -\frac{3}{2}, \frac{1}{2}$$

$$J_2^2 |j_2, m_2\rangle = 6 \hbar^2 |j_2, m_2\rangle = j_2(j_2 + 1) \hbar^2 |j_2, m_2\rangle$$

$$j_2^2 + j_2 - 6 = 0 \Rightarrow j_2 = -3, 2$$

So we have $j_1 = \frac{1}{2}, j_2 = 2$. Using the Clebsch-Gordan Coefficient table, and looking for

$|j, m\rangle \Rightarrow |j_2, j_1, m_2, m_1\rangle$, we get a set of simultaneous normalized eigenvectors,

$$\left| \frac{5}{2}, \frac{5}{2} \right\rangle = \left| 2, \frac{1}{2}, 2, \frac{1}{2} \right\rangle$$

$$\left| \frac{5}{2}, \frac{3}{2} \right\rangle = \frac{1}{\sqrt{5}} \left| 2, \frac{1}{2}, 2, -\frac{1}{2} \right\rangle + \frac{2}{\sqrt{5}} \left| 2, \frac{1}{2}, 1, \frac{1}{2} \right\rangle$$

$$\left| \frac{3}{2}, \frac{3}{2} \right\rangle = \frac{2}{\sqrt{5}} \left| 2, \frac{1}{2}, 2, -\frac{1}{2} \right\rangle - \frac{1}{\sqrt{5}} \left| 2, \frac{1}{2}, 1, \frac{1}{2} \right\rangle$$

$$\left| \frac{5}{2}, \frac{1}{2} \right\rangle = \sqrt{\frac{2}{5}} \left| 2, \frac{1}{2}, 1, -\frac{1}{2} \right\rangle + \sqrt{\frac{3}{5}} \left| 2, \frac{1}{2}, 0, \frac{1}{2} \right\rangle$$

$$\left| \frac{3}{2}, \frac{1}{2} \right\rangle = \sqrt{\frac{3}{5}} \left| 2, \frac{1}{2}, 1, -\frac{1}{2} \right\rangle - \sqrt{\frac{2}{5}} \left| 2, \frac{1}{2}, 0, \frac{1}{2} \right\rangle$$

$$\left| \frac{5}{2}, -\frac{1}{2} \right\rangle = \sqrt{\frac{3}{5}} \left| 2, \frac{1}{2}, 0, -\frac{1}{2} \right\rangle + \sqrt{\frac{2}{5}} \left| 2, \frac{1}{2}, -1, \frac{1}{2} \right\rangle$$

$$\left| \frac{3}{2}, -\frac{1}{2} \right\rangle = \sqrt{\frac{2}{5}} \left| 2, \frac{1}{2}, 0, -\frac{1}{2} \right\rangle - \sqrt{\frac{3}{5}} \left| 2, \frac{1}{2}, -1, \frac{1}{2} \right\rangle$$

$$\left| \frac{5}{2}, -\frac{3}{2} \right\rangle = \frac{2}{\sqrt{5}} \left| 2, \frac{1}{2}, -1, -\frac{1}{2} \right\rangle + \frac{1}{\sqrt{5}} \left| 2, \frac{1}{2}, -2, \frac{1}{2} \right\rangle$$

$$\left| \frac{3}{2}, -\frac{3}{2} \right\rangle = \frac{1}{\sqrt{5}} \left| 2, \frac{1}{2}, -1, -\frac{1}{2} \right\rangle - \frac{2}{\sqrt{5}} \left| 2, \frac{1}{2}, -2, \frac{1}{2} \right\rangle$$

$$\left| \frac{5}{2}, -\frac{5}{2} \right\rangle = \left| 2, \frac{1}{2}, -2, -\frac{1}{2} \right\rangle$$

Question 4(a)

Let c_m be the number of states $|j, m\rangle$ for a fixed $m = m_1 + m_2$, and d_m be the number of states $|j, m\rangle$ for a fixed m and j . With $|m| \leq j$, and assume $m \geq 0$,

$$c_m = \sum_{m \geq 0} d_m = d_m + d_{m+1} + \dots$$

$$c_{m+1} = d_{m+1} + d_{m+2} + \dots$$

$$c_m - c_{m+1} = d_m$$

m	m_1	m_2	c_m
$j_1 + j_2$	j_1	j_2	1
$j_1 + j_2 - 1$	$j_1 - 1$ j_1	j_2 $j_2 - 1$	2
$j_1 + j_2 - 2$	$j_1 - 2$ $j_1 - 1$ j_1	j_2 $j_2 - 1$ $j_2 - 2$	3
...
$j_1 + j_2 - n$			$n + 1$

So if you know c_m , we can get d_m . As in the case of a simple harmonic oscillator, we know that $d_j = 1$ if j is allowed, and $d_j = 0$ if j is not allowed. The value of n to terminate the counting is $-2j_2$ if $j_2 < j_1$, or $-2j_1$ if $j_1 < j_2$.

$$\therefore j = j_1 + j_2, j_1 + j_2 - 1, j_1 + j_2 - 2, \dots, |j_1 - j_2| \quad [\text{shown}]$$

Question 4(b)

$$E = \frac{\hbar^2 \pi^2}{2mb^2} (k^2 + l^2 + m^2 + n^2)$$

(i) For distinguishable identical particles,

$$\psi = \frac{4}{b^2} \sin\left(\frac{k\pi x_1}{b}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_3}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right)$$

Ground state, $E = \frac{2\hbar^2 \pi^2}{mb^2}$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

1st excited state, $E = \frac{7\hbar^2 \pi^2}{2mb^2}$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{2\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{2\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{2\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

