

## PC5215, Numerical Recipes with Applications

### Lab 4, due Friday 13 November 2020

1. In this last lab, we consider solving one-electron time-dependent Schrödinger equation in one dimension scattering over a potential barrier. We send an electron from the left side of the barrier, and ask what is the probability  $T(E)$  (transmission probability) that the electron passes through the barrier, as a function of incoming electron energy  $E$ . The equation is

$$i\hbar \frac{\partial \Psi}{\partial t} = H\Psi, \quad H = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x).$$

We choose two forms of the potential

$$V_1(x) = \begin{cases} 0, & x < 0, \\ V_0, & 0 \leq x < a, \\ 0, & x \geq a. \end{cases} \quad V_2(x) = \begin{cases} 0, & x < 0, \\ \frac{4V_0}{a^2}(a-x)x, & 0 \leq x < a. \\ 0, & x \geq a. \end{cases}$$

For numerical computation, we take  $V_0 = 1$  eV, and take the distance of the barrier  $a$  such

that  $\frac{\hbar^2}{2ma^2} = \frac{1}{100}$  eV.

- a. First, we solve the problem with the rectangular potential  $V_1$  analytically as a check for part (b) and (c) using numerical methods. To do this, use the plane-wave as a trial solution  $\Psi(x,t) = ce^{i(\pm px - Et)/\hbar}$  and match the boundary conditions (the function and its first derivative should be continuous) at 0 and  $a$  to find the whole solution. Determine the transmission probability by the absolute value square of the outgoing wave ( $T = |c|^2$ ) if the incoming wave has amplitude 1. Note that  $T + R = 1$ , where  $R$  is the probability that the particle is reflected back to the left. Plot  $T$  as a function of  $E$  from 0 to 2 eV.
- b. Next, we solve the same problem as in (a) numerically using “wave packet”. The idea is that we send a wavepacket from the left side with energy centered around  $E = \frac{p_0^2}{2m}$  and evolve the wavepacket in time for sufficiently long-time and then ask what is the total probability that the particle is on the right side. We use the following form for the Gaussian wavepacket with position centered around  $x_0$  and momentum center around  $p_0$ , as the initial condition to the time-dependent Schrödinger equation:

$$\Psi(x) \propto \exp\left(-\frac{1}{2\sigma^2}(x-x_0)^2 + \frac{i}{\hbar}p_0(x-x_0)\right).$$

To evolve the wavepacket, we use the “operator splitting method”. First, note that formally, the solution can be found by  $\Psi(t+\Delta t) = e^{-\frac{i}{\hbar}H\Delta t}\Psi(t)$ . We then use Trotter-Suzuki formula

$$e^{A+B} \approx e^A e^B$$

where  $A$  and  $B$  are arbitrary operators but small of order  $\Delta t$  (a more accurate one is  $e^{A+B} \approx e^{A/2} e^B e^{A/2}$ ). We take the kinetic energy term as  $A$  and potential energy

term as  $B$ . In coordinate representation,  $e^B\Psi = e^{-\frac{i}{\hbar}\Delta t V(x)}\Psi(x)$ , which is very easy computationally, but the kinetic energy part is hard because it is a differential operator. However, it will be easy if we work in momentum representation. Then

it becomes  $e^A\Psi = e^{-\frac{i}{\hbar}\Delta t \frac{p^2}{2m}}\Psi(p)$ . The relationship between coordinate and momentum representation is obtained by Fourier transform (using the standard convention in quantum mechanics):

$$\Psi(p) = \langle p | \Psi \rangle = \int_{-\infty}^{+\infty} dx \langle p | x \rangle \langle x | \Psi \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} \Psi(x) e^{-i\frac{px}{\hbar}} dx.$$

To implement a Schrödinger equation solver on computer, (i) discretize the coordinate  $x$  and momentum  $p$  appropriately so that FFT can be conveniently performed. You could use `numpy.fft` under Python. (ii) multiply the associated phase factor for the Trotter-Suzuki decomposition to evolve  $\Psi$  in time, changing the representations whenever it is needed. Compare your answers for  $T(E)$  with the exact one find in part (a). Pay attention to the unspecified parameters such as  $x_0$ ,  $\sigma$ ,  $\Delta t$ , discretization steps for  $x$  and  $p$  and number of sampling point  $N$  in FFT, so that your answers are reasonably accurate. [Hint:  $\sigma$  should be large otherwise it is not close to a plane wave;  $N$  should be of the order  $10^4$ . You better make plots to see where your wavepacket is after a contain time step.]

- c. Repeat the same calculation as in (b) with the parabolic potential  $V_2$ . Compare and discuss the result.
- d. Alternative, possibly better methods not using FFT: Crank-Nicholson method. This will be a bonus question if some of you wish to try and compare with FFT/Trotter-Suzuki method in part b.